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## Absence of magnetic thermal conductivity in the quantum spin liquid candidate YbMgGaO₄

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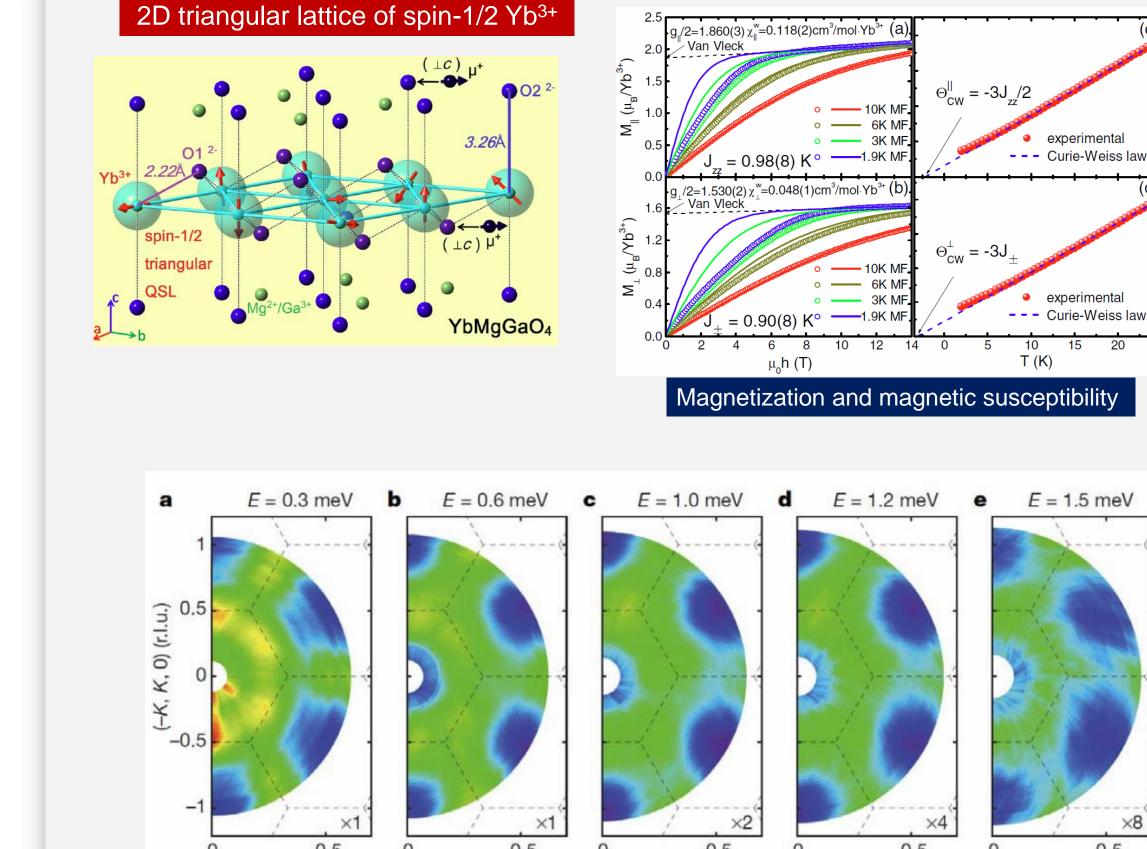
We present the ultralow-temperature specific heat and thermal conductivity measurements on single crystals of YbMgGaO<sub>4</sub>, which was recently argued to be a promising candidate for a quantum spin liquid (QSL). In a zero magnetic field, a large magnetic contribution of specific heat is observed, and exhibits a power law temperature dependence ( $C_m \sim T^{0.74}$ ). On the contrary, we do not observe any significant contribution of thermal conductivity from magnetic excitations. In magnetic fields  $H \ge 6 T$ , the exponential T dependence of  $C_m$  and the enhanced thermal conductivity indicate a magnon gap of the fully polarized state. The absence of magnetic thermal conductivity at the zero field in this QSL candidate puts a strong constraint on the theories of its ground state.

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(a.u.)

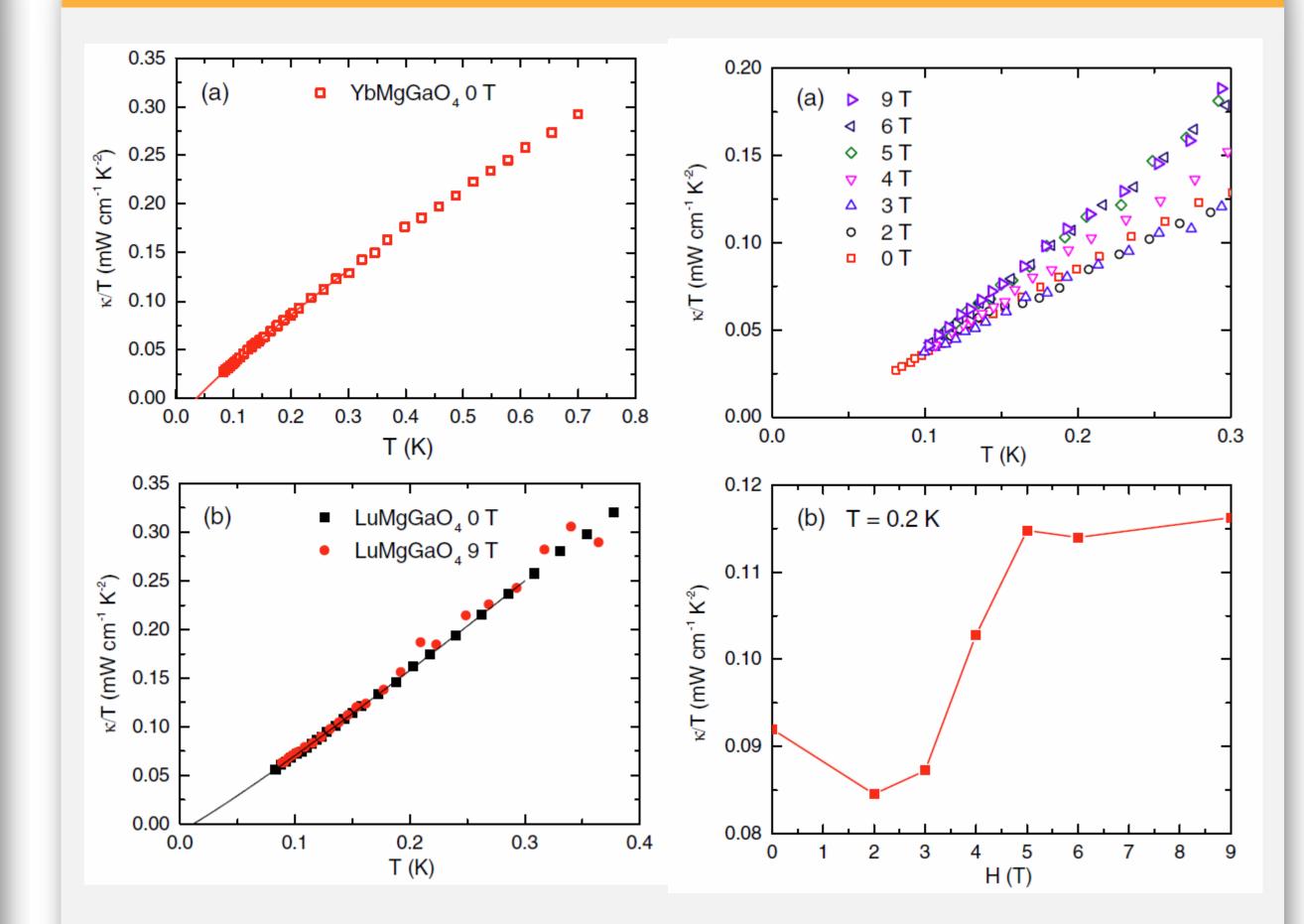
### **Thermal transport properties**

## YbMgGaO₄: a QSL candidate

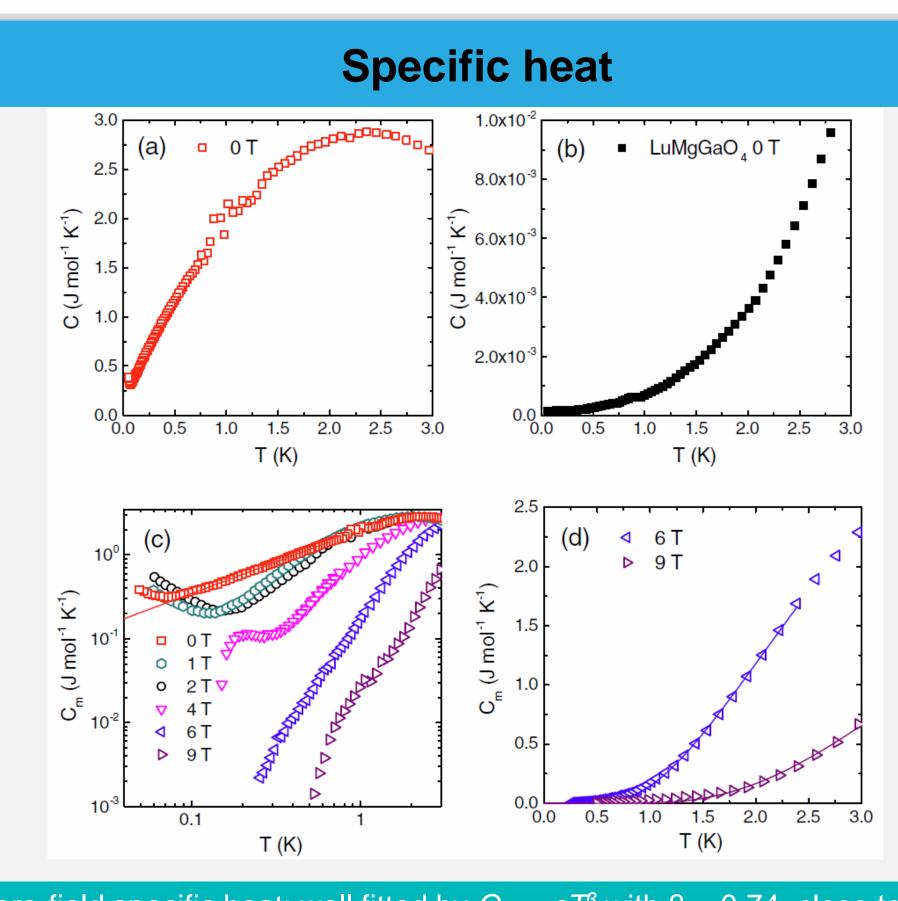


 $\times 8$ 0.5 0.5 0.5 0.5 0.5 0 (H, H, 0) (r.l.u.) (H, H, 0) (r.l.u.) (H, H, 0) (r.l.u.) (H, H, 0) (r.l.u.) (H, H, 0) (r.l.u.)

A broad continuum of spin excitations, which is a hallmark of the QSL state, was observed in neutron scattering measurements



At zero field,  $\kappa_0/T$  is negative, and the power of the temperature dependence is abnormally lower



Zero-field specific heat: well fitted by  $C_m = cT^{\beta}$  with  $\beta = 0.74$ , close to 2/3.

A gapless U(1) QSL with a spinon Fermi surface according to Ref. [5]

Magnetic field rapidly suppresses the  $C_m$ . The temperature dependence of the  $C_m$ 

than 2. The  $\kappa/T$  first decreases slightly for H < 2 T, then there is a sharp increase between 2 and 5 T, and it finally saturates for H > 5 T. While the magnetic state of YbMgGaO<sub>4</sub> in the intermediate fields (0 < H < 5T) is rather complex, it simply tends to become a fully polarized state for H > 5Tat such low temperatures. In the fully-polarized state with a magnon gap of several Kelvins, there are almost no magnetic excitations to scatter phonons below 0.3 K; therefore, the  $\kappa$  of YbMgGaO<sub>4</sub> at H > 5 T is purely contributed by phonons. At lower fields, it is the additional scattering of phonons by magnetic excitations that suppresses the κ and gives the abnormal temperature dependence of  $\kappa$  and the unphysical negative  $\kappa_0/T$  for YbMgGaO<sub>4</sub>.

## Discussions

> The absence of magnetic thermal conductivity at zero field means that either (a) the presumed gapless spinons do not exist in YbMgGaO<sub>4</sub> and the large  $C_m$ has some other magnetic origin;

or (b) the gapless spinons do exist but for some reason they do not conduct heat significantly in YbMgGaO<sub>4</sub>. One possible mechanism of the spinon localization may be the disorder of Mg<sup>2+</sup>-Ga<sup>3+</sup> sites (random occupation) in the double layers of Mg/GaO<sub>5</sub> triangular bipyramids.

References

#### (Only major ones listed):

[2] M. Yamashita et al., Science **328**, 1246 (2010). [3] Y. S. Li et al., Sci. Rep. 5, 16419 (2015).

[1] M. Yamashita et al., Nat. Phys. 5, 44 (2009). [4] Y. S. Li et al., Phys. Rev. Lett. 115, 167203 (2015). [5] Y. Shen *et al.*, Nature **540**, 559 (2016). [6] J. A. M. Paddison *et al.*, Nat. Phys. **13**, 117 (2017).

#### gradually turns into an exponential one, which is attributed to the magnons with a gap in a

fully polarized state.

#### Y. Xu *et al.*, Physical Review Letters 117, 267202 (2016).