## Broken time-reversal symmetry in superconducting partially-filled skutterudite $Pr_{1-\delta}Pt_4Ge_{12}$

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## Background

 $\succ$  The superconducting gap structure in PrPt<sub>4</sub>Ge<sub>12</sub> is controversial: Point node?[1]? BCS[2]? Multiband[3][4]?

 $\succ$  It is puzzling that time reversal symmetry (TRS) breaks at temperature  $T_m < T_c$ .

## **Motivations**

A study of an imperfectly filled skutterudite  $Pr_{1-\delta}Pt_4Ge_{12}$  to investigate;

(i) whether the cage-forming structure and superconductivity survive with partially filled <sup>141</sup>Pr nuclei; (ii) the effect of the insufficient filled <sup>141</sup>Pr nuclei on TRS breaking in PrPt<sub>4</sub>Ge<sub>12</sub>

(iii) the gap symmetry of  $Pr_{1-\delta}Pt_4Ge_{12}$  and implications of the superconducting order parameter of  $PrPt_4Ge_{12}$ .





(a) Magnetic susceptibility of the sample (b) Specific FIG. 1 heat data of the sample displayed as  $C_p/T$  versus  $T^2$ .

> Superconductivity is observed below  $T_c = 7.80$  K, the same as PrPt<sub>4</sub>Ge<sub>12</sub>.





(a) Zero-field  $\mu$ SR time spectra at 12 K (red cir-FIG. 3 cles) and 0.6 K (blue squares) for the sample. A background signal has been subtracted from the data. The corresponding solid lines are fits according to Eq. (3), where  $\lambda$  was fixed at 0  $\mu s^{-1}$ . (b) Temperature dependence of the muon spin relaxation rates  $\sigma$  (red circles) and  $\Lambda$  (blue triangles).  $\sigma$  was derived from the fitting of Eq. (3) with  $\Lambda$  fixed at 0.08  $\mu s^{-1}$ . The red curve is the fit of Eq. (5). The blue line denotes the average of  $\Lambda$  data from 0 to 13 K.

$$P_{\mu}(t) = G_{z}^{K-T}(\sigma, \lambda, t) \exp(-\Lambda t).$$
(3)

Temperature dependence of the electronic specific FIG. 2 heat coefficient of  $Pr_{1-\delta}Pt_4Ge_{12}$  at zero field. The curves represent the fits using six different gap models. Inset: the critical fields  $H_{c2}$  derived from the midpoints of the jump in  $C_p/T$ . The green curve is the fit of Eq. (1) to data.

 $\mu_0 H_{c2}(T) = \mu_0 H_{c2}(0) \frac{1 - t^2}{1 + t^2}, \quad (1)$ 

> Temperature dependence of both the upper critical field and the electronic specific heat can be described in terms of a two-gap model: evidence of multi-band superconductivity.

$$G_{z}^{K-T}(\sigma,\lambda,t) = \frac{1}{3} + \frac{2}{3}(1 - \sigma^{2}t^{2} - \lambda t)\exp(-\frac{1}{2}\sigma^{2}t^{2} - \lambda t)(4)$$
$$\sigma_{n}, T > T_{c}, (5)$$
$$[\sigma_{n}^{2} + \sigma_{e}(T)^{2}]^{1/2}, T < T_{c}, (5)$$

 $\succ \sigma_{e}(0)/\gamma_{\mu}$  in  $Pr_{1-\delta}Pt_{4}Ge_{12}$  is 0.077(4) mT, nearly half of that (0.141 mT) in  $PrPt_4Ge_{12}$ .

## Conclusions

- $\succ$  The [Pt<sub>4</sub>Ge<sub>12</sub>] cage-forming structure survives and superconductivity is observed below  $T_c$  = 7.80 K.
- $\succ$  The temperature dependence of  $Hc_2$  and the electronic specific heat are well described by the twoband model.
- > The onset of broken TRS is observed at  $T_m < T_c$ , possibly due to the appearance of a second phase, while no obvious specific jump is observed around  $T_m$ .  $\succ \sigma_{e}(0)/\gamma_{\mu}$  in  $Pr_{1-\delta}Pt_{4}Ge_{12}$  is half of that in  $PrPt_{4}Ge_{12}$ , indicating that the<sup>141</sup>Pr nuclei or Pr-Pr intersite interactions are responsible for broken TRS.

