

Unusual slow magnetic fluctuations in $\text{Sr}_2\text{Ir}_{1-x}\text{Rh}_x\text{O}_4$ studied by muon spin relaxation (μSR)

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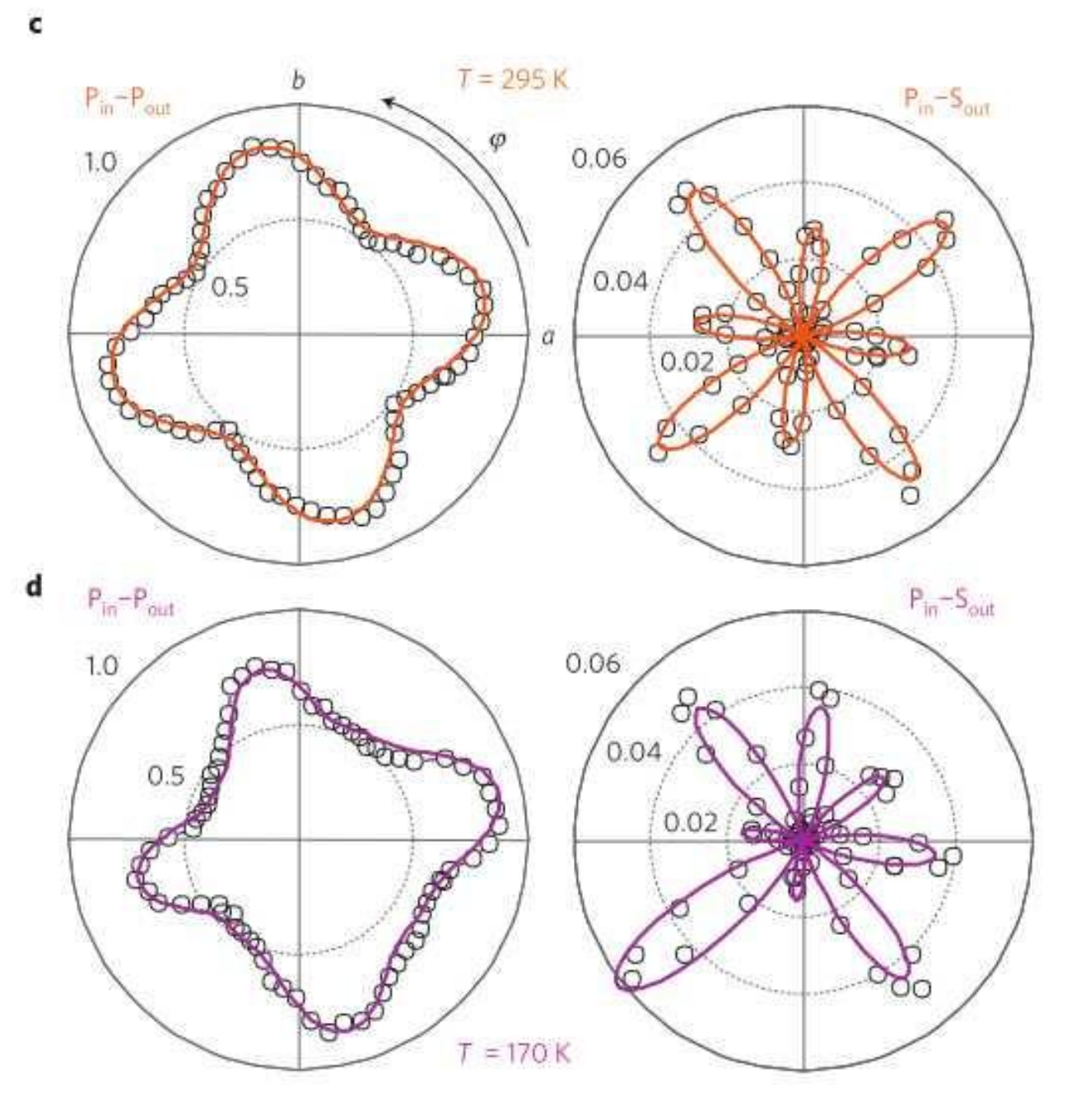
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Motivations : The hidden order phase is similar to pseudo-gap state in cuprates

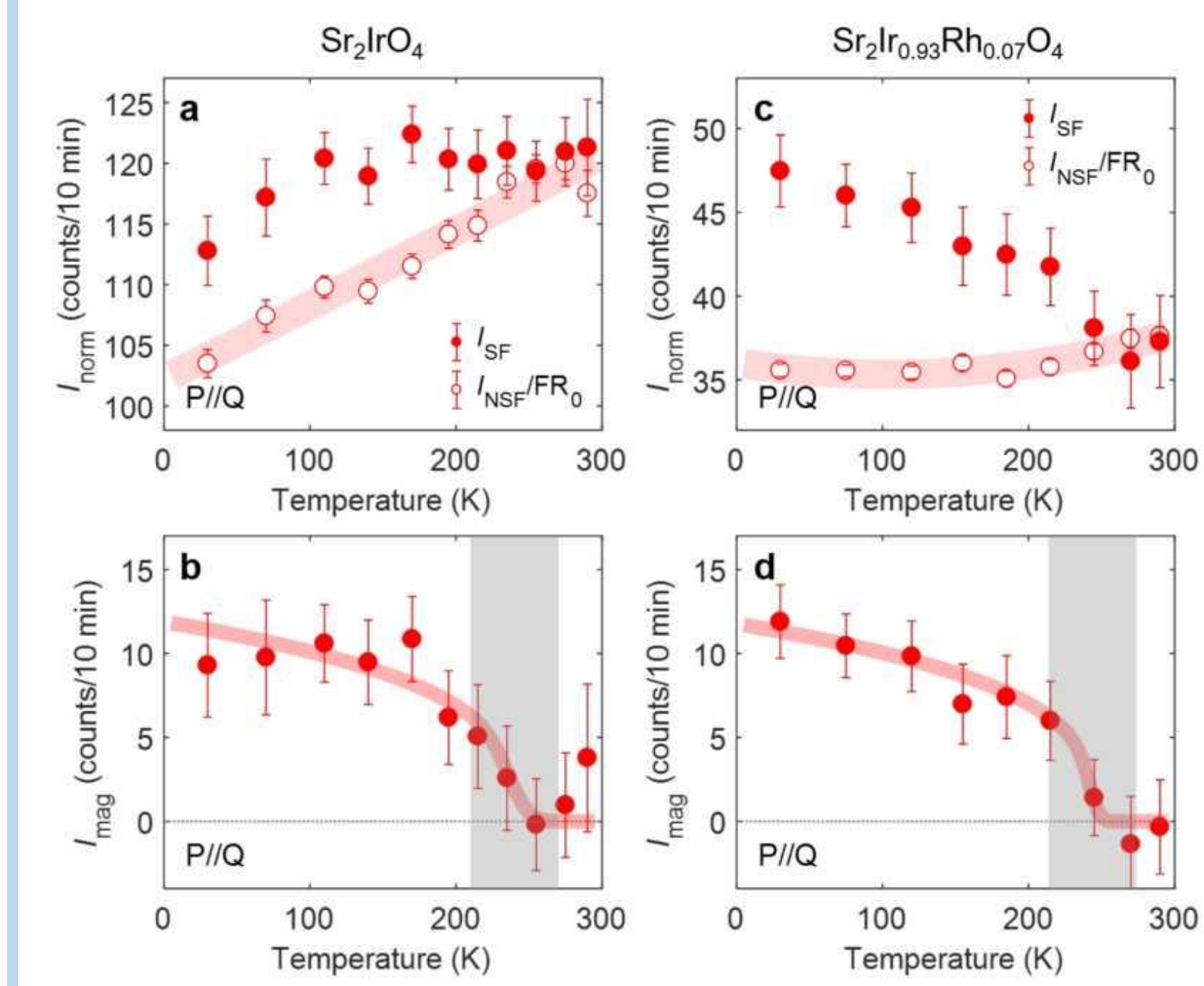
Rotational symmetry breaking

Time reversal symmetry breaking

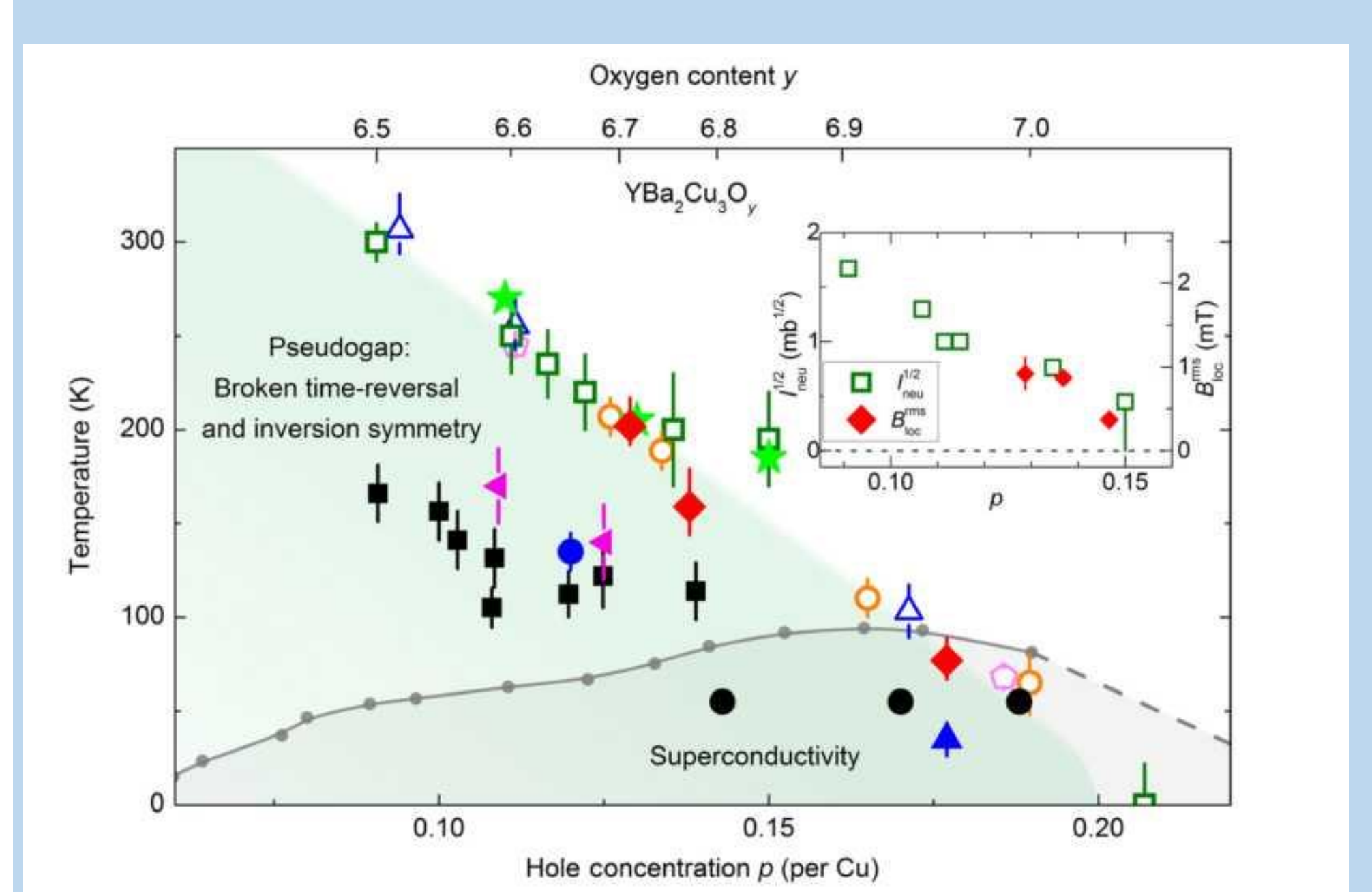
Pseudo-gap state in cuprates



Zhao, L. *et al* Nature_Physics (2015)



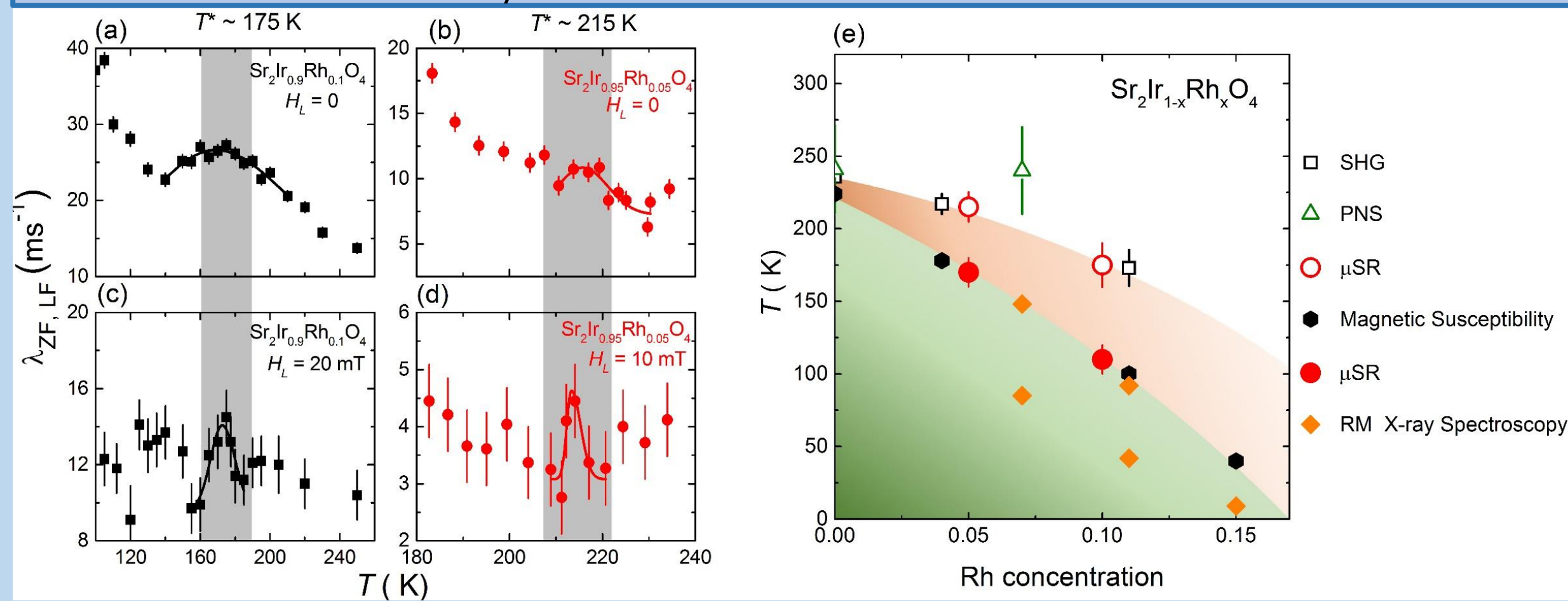
Jeong, J. *et al*. NATURE COMMUNICATIONS (2017)



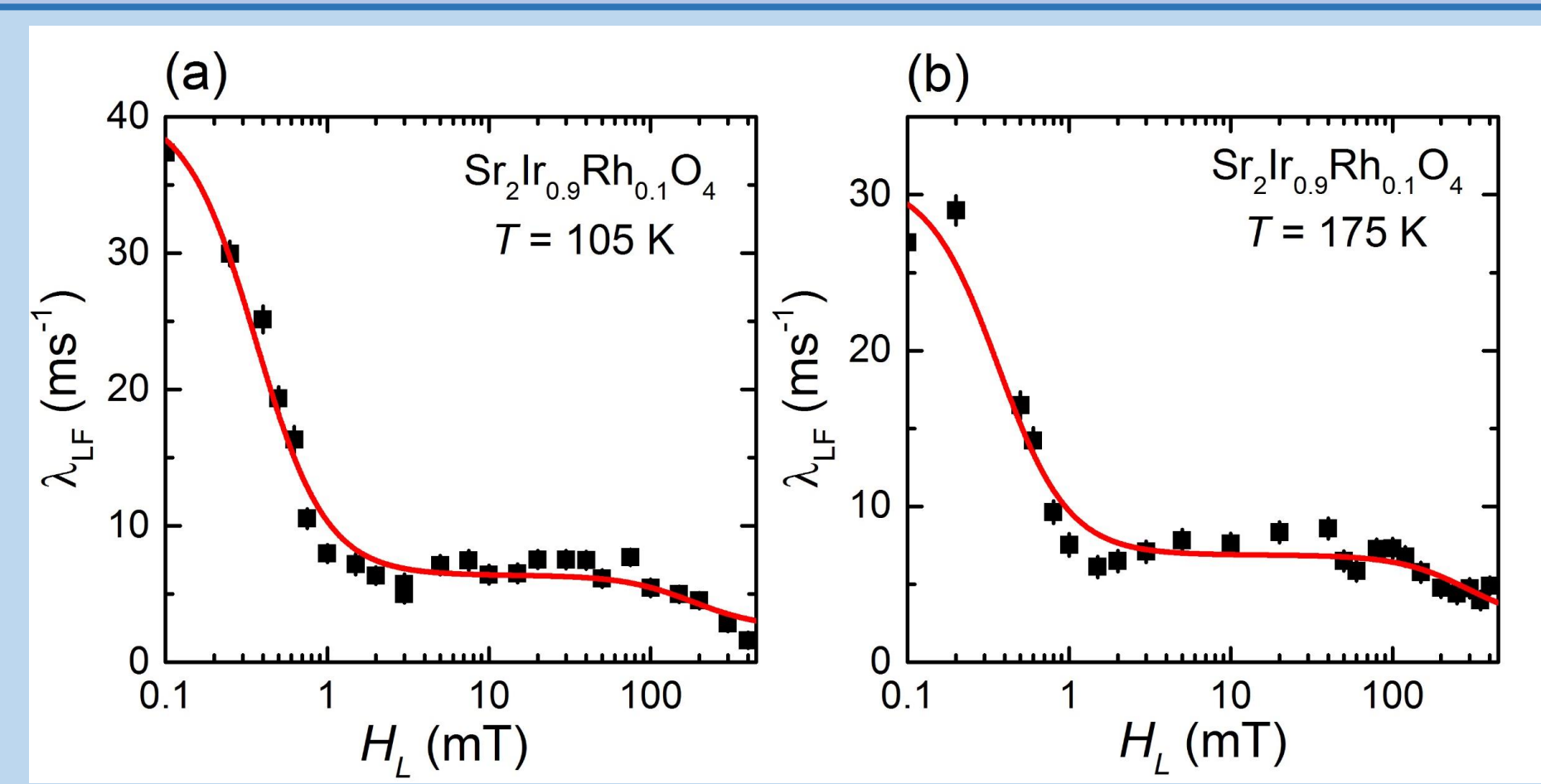
Zhang, Ding *et al*. Science Advances (2018)

ZF & LF- μSR Results

$\lambda_{ZF/LF}(T)$: The maxima at T^*



$\lambda_{LF}(H_L)$: Two slow magnetic fluctuations

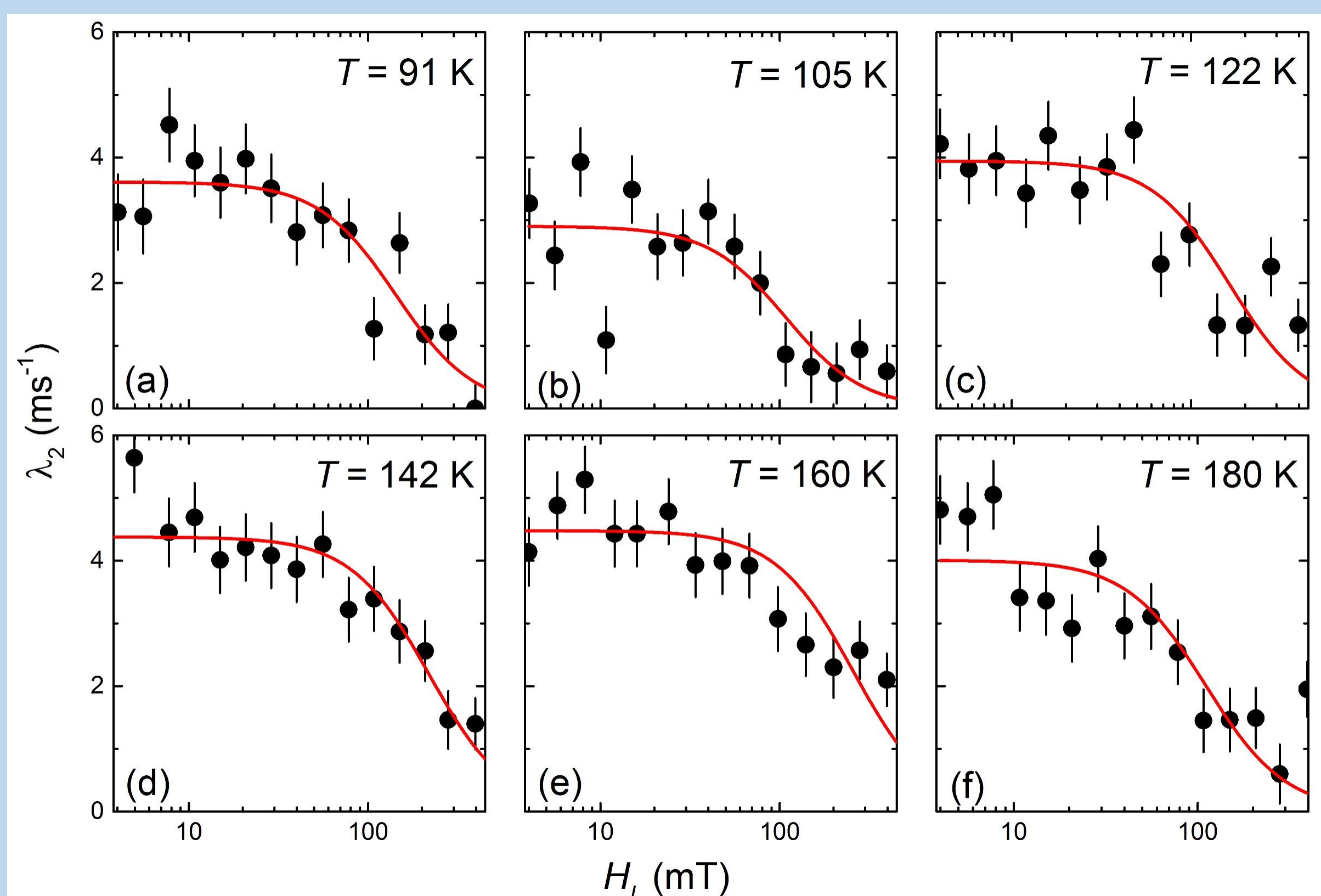


Redfield Relation: The function to describe field dependence of relaxation rate

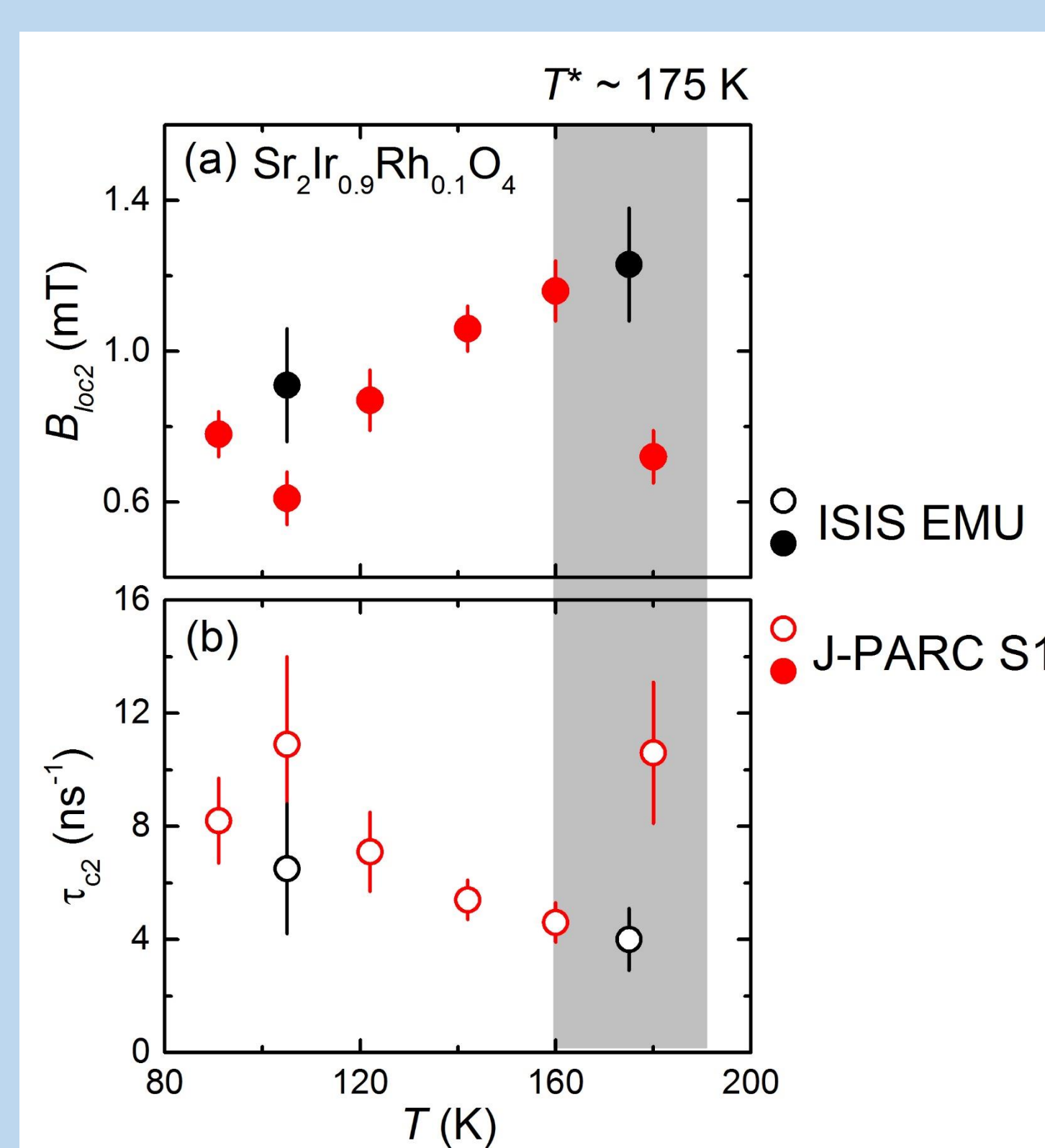
$$\lambda_{LF}^{(1)} = \frac{2(\gamma_{\mu} B_{loc}^{(1)})^2 \tau_c^{(1)}}{1 + (\tau_c^{(1)} \gamma_{\mu} H_L)^2}, \quad \lambda_{LF}^{(2)} = \frac{2(\gamma_{\mu} B_{loc}^{(2)})^2 \tau_c^{(2)}}{1 + (\tau_c^{(2)} \gamma_{\mu} H_L)^2}$$

$\lambda_{ZF/LF} = \lambda_{LF}^{(1)} + \lambda_{LF}^{(2)}$, $\lambda_{LF}^{(1)}$ represents the field dependence of λ_{LF} at low fields, corresponding to **nuclear dipolar field**; $\lambda_{LF}^{(2)}$ represents the high field part, corresponding to the **hidden magnetic order**.

$\lambda_{LF}(H_L)$ at different temperatures



B_{loc} and τ_c versus T



- Positive slope of $B_{loc}^{(2)}(T)$
- Negative slope of $\tau_c^{(2)}$
- Critical slowing down at T^*



Preliminary conclusions

- The temperature of maxima in $\lambda_{ZF}(T)$ are consistent with its measured by symmetry breaking probe.
- Our LF results indicate that there are two different magnetic fluctuations with different fluctuating rate in $\text{Sr}_2\text{Ir}_{0.9}\text{Rh}_{0.1}\text{O}_4$.
- The maxima we observed in ZF is caused by faster magnetic fluctuation.
- The maxima in $\lambda_{ZF}(T)$ indicates a critical slowing down

