

Pressure-Tuned Topological Phases in Kagome Magnets (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub>

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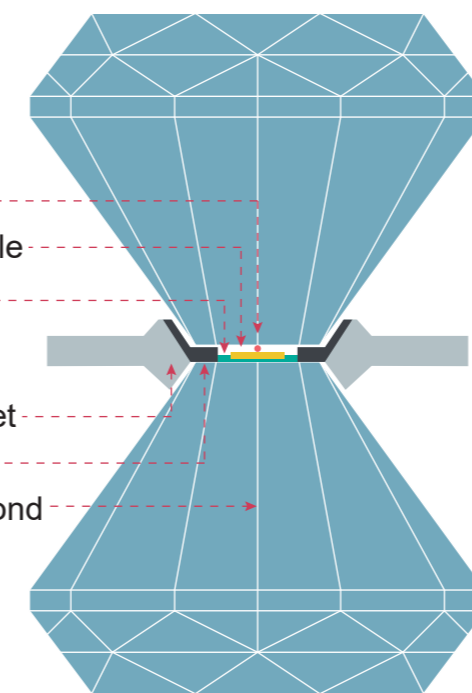
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## Introduction

As one of the origins of novel phenomena encompassing quantum magnetism, a kagome lattice has played a key role for realizing and tuning novel electronic states, owing to its unique lattice geometry and band structure.

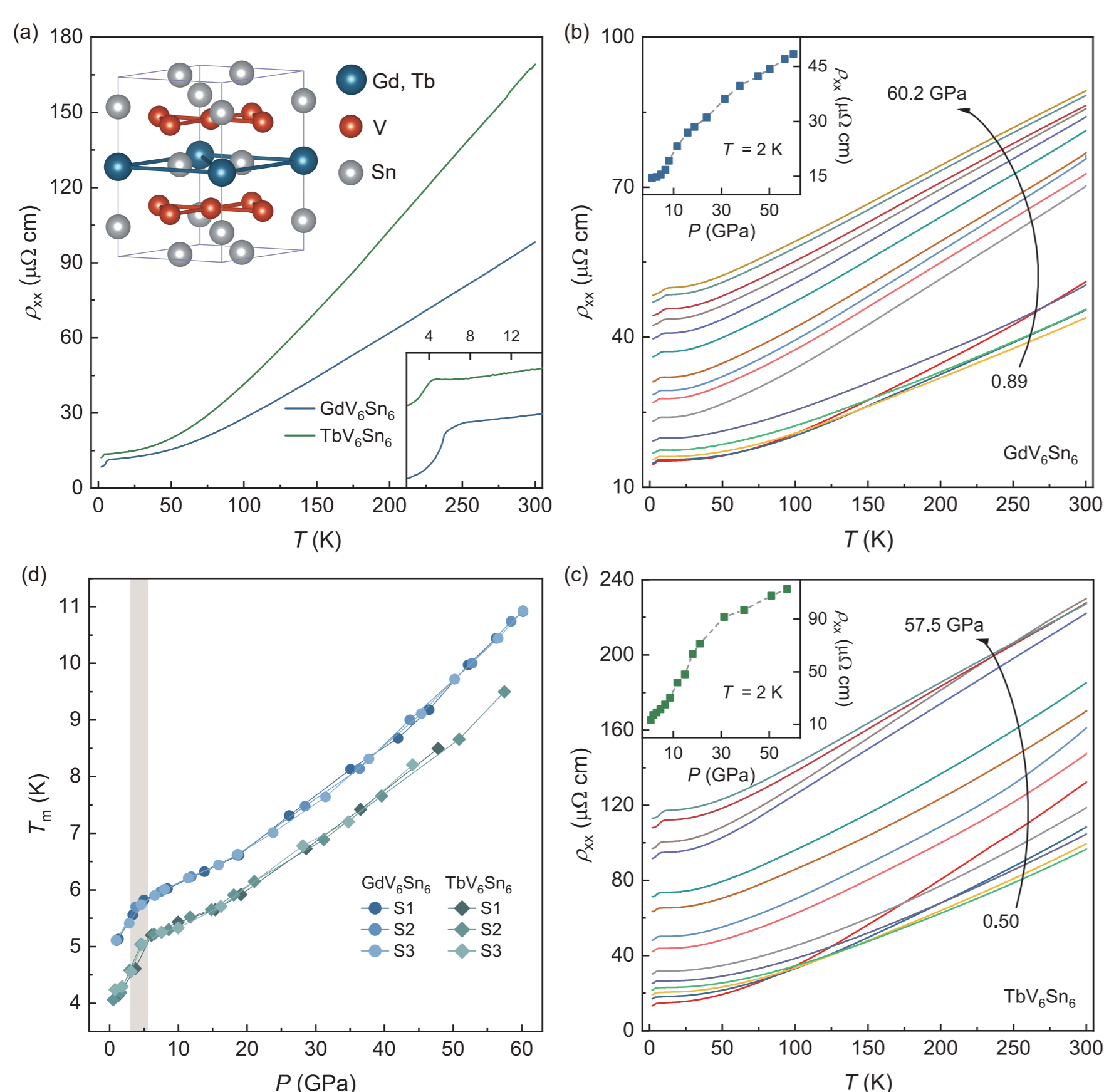
RV<sub>6</sub>Sn<sub>6</sub> (R = rare-earth element) is a newly discovered kagome family, which is isostructural to TbMn<sub>6</sub>Sn<sub>6</sub> structural prototype but the kagome layer composed of V atoms is nonmagnetic. For magnetic R-triangular lattice (R = Gd-Tm), the segregated magnetic layer not only permits a direct tunability for nonmagnetic kagome layer below the magnetic transition temperature  $T_m$ , but also obviates the complex magnetic configurations in RMn<sub>6</sub>Sn<sub>6</sub> family because of a pristine Mn-kagome lattice.

## Method

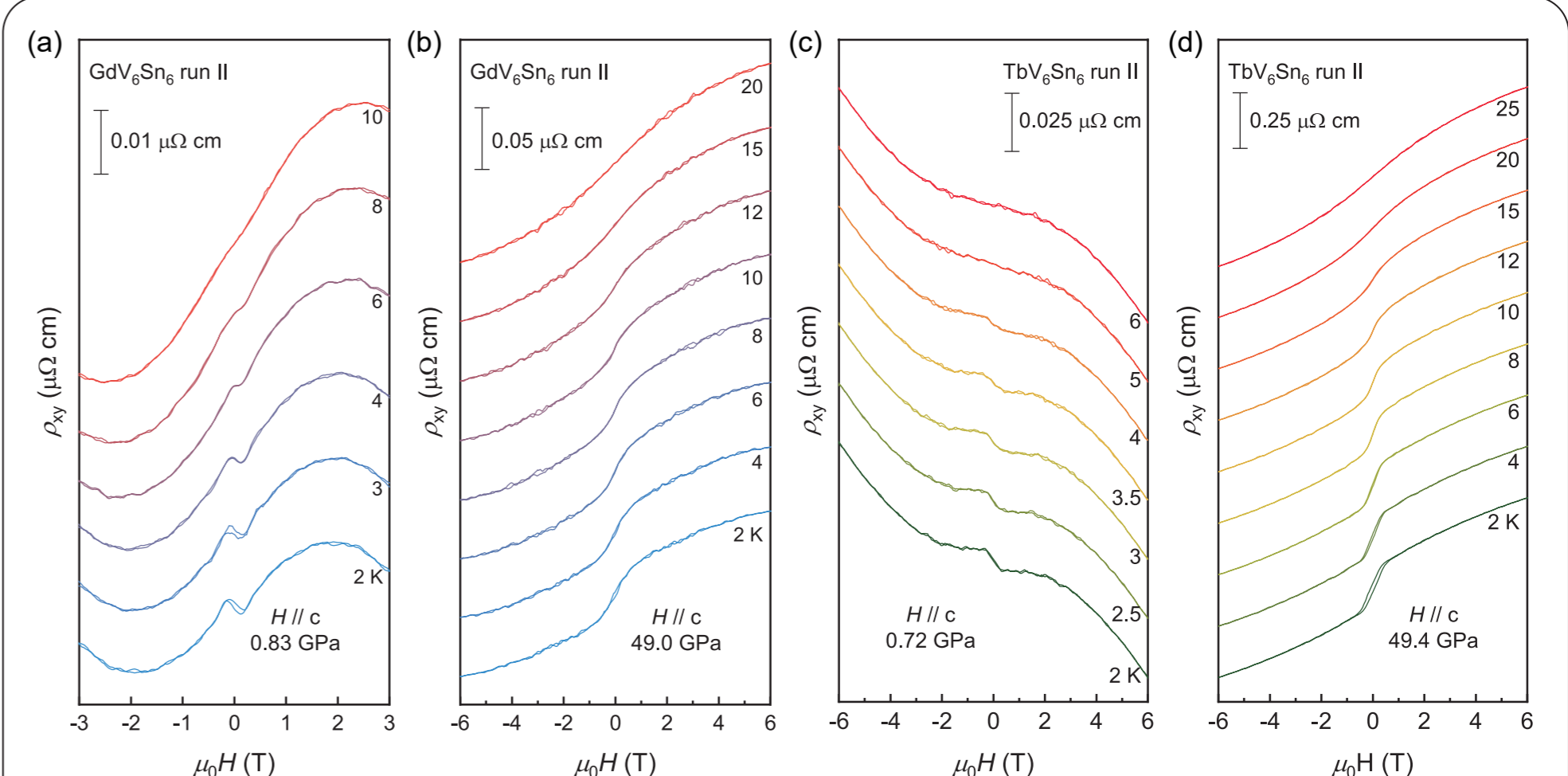


High-pressure electric transport measurements were performed in a diamond anvil cell (DAC) made of BeCu alloy with NaCl powder as the pressure transmitting medium. A mixture of c-BN powder and epoxy is adopted as the insulating layer. The pressures are determined by the ruby luminescence method.

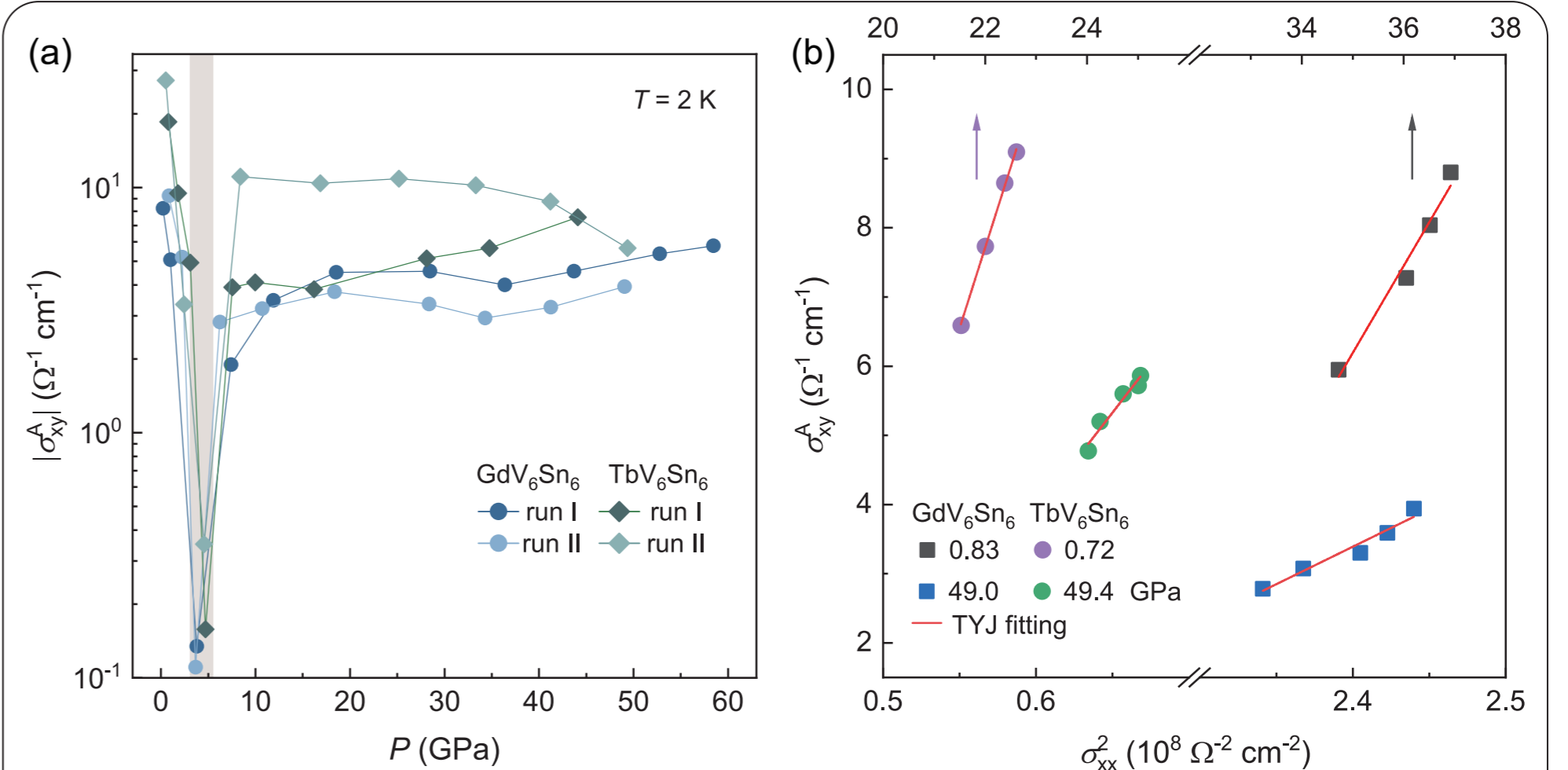
The four-probe method and Van Der Pauw method are applied for the transverse magnetoresistance and Hall resistance measurements, respectively.

 Longitudinal resistivity of (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub> at different pressures


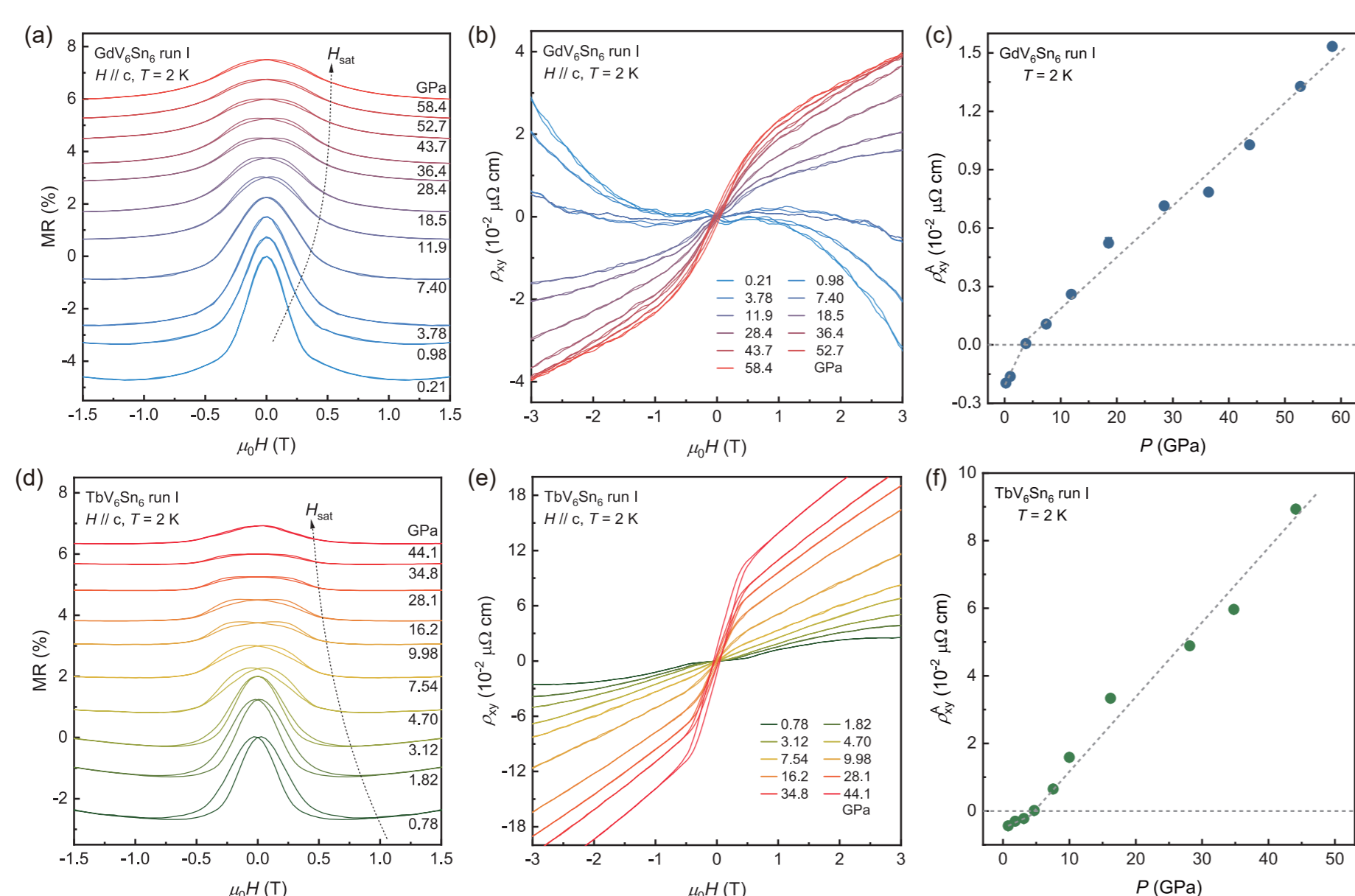
(a) Longitudinal resistivity  $\rho_{xx}$  as a function of temperature  $T$  at ambient pressure. Inset: crystal structure of (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub> and closeup of  $\rho_{xx}$  at  $T = 2$  K. (b), (c)  $\rho_{xx}(T)$  curves under various pressures. Inset: resistivity under different pressure at  $T = 2$  K. (d) Pressure dependence of magnetic transition temperature  $T_m$ .

 Temperature-dependent AHE of (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub> under pressure


(a)-(d)  $\rho_{xy}$  versus  $\mu_0 H$  at various temperatures. Both of samples exhibit a positive anomalous Hall effect (AHE) under high pressure, while low-pressure AHE responds oppositely. The negative AHE is robust although the low-pressure  $\rho_{xy}$  of GdV<sub>6</sub>Sn<sub>6</sub> represents another behavior against bulk or other batches in the range of  $\mu_0 H = \pm 2$  T.

 Phase diagram and TYJ model fittings for (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub>


(a) Phase diagram of GdV<sub>6</sub>Sn<sub>6</sub> and TbV<sub>6</sub>Sn<sub>6</sub> with pressure. Beige shaded area highlights the anomaly of anomalous Hall conductivity  $\sigma_{xy}^A$  at  $T = 2$  K crossing the critical pressures. (b) Fittings of  $\sigma_{xx}^2$  dependence of  $\sigma_{xy}^A$  using TYJ model (PRL 103, 087206 (2009)) for GdV<sub>6</sub>Sn<sub>6</sub> and TbV<sub>6</sub>Sn<sub>6</sub>. All of well linear relationships reflect the pressure-induced AHE can mainly result from the intrinsic effect of the Berry curvature.

 Pressure tuned magnetotransports of (Gd, Tb)V<sub>6</sub>Sn<sub>6</sub>


(a), (d) Field dependence of transverse magnetoresistivity (MR) of GdV<sub>6</sub>Sn<sub>6</sub> and TbV<sub>6</sub>Sn<sub>6</sub> under different pressure at  $T = 2$  K. (b), (e) Hall resistivity  $\rho_{xy}$  under varying pressures at  $T = 2$  K. (c), (f) Anomalous Hall resistivity  $\rho_{xy}^A$  as a function of pressure, which is derived from (b) and (e), respectively.

## Conclusions

- In GdV<sub>6</sub>Sn<sub>6</sub> and TbV<sub>6</sub>Sn<sub>6</sub>, we observe two pressure-induced anomalous Hall effect (AHE) features: an unusual suppression (below critical pressure) and a progressive enhancement (above critical pressure).
- Based on the TYJ model, both AHE features are mainly attributed to the intrinsic contribution of the Berry curvature. The AHE-weakened and AHE-enhanced regimes indicate that two distinct modulations of topological bands before and after the critical pressure.