

A new type of stable borophene with flat-band-induced magnetism



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Abstract

Based on first-principles calculations, we propose a new type of thermally and dynamically stable magnetic borophene (B_{11}) with a tetragonal lattice. The magnetism is found coming from spin polarization of one bonding flat band located at the Fermi level. The strong interactions between the p_z orbitals of the B atoms in the structural unit lead to the formation of the flat bands. One tight binding model is built to comprehend the magnetic mechanism. Biaxial tensile strain (2.0%) can trigger a phase transition from a semimetal to a semiconductor in the B_{11} monolayer. Our work provides a new route for designing and achieving two-dimensional magnetic materials with light elements.

Geometry structure and its stability

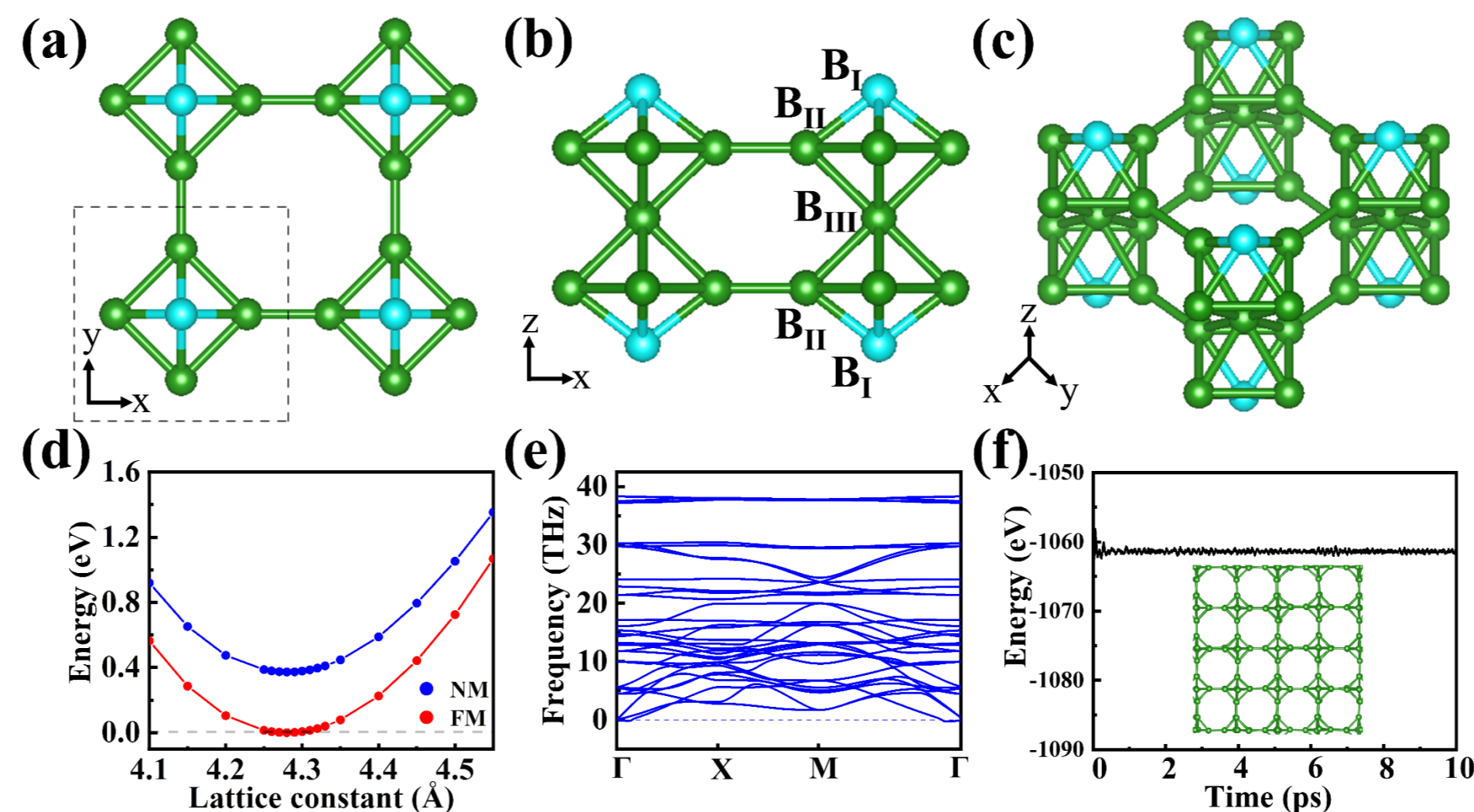


FIG. 1: (a) Top view of the B_{11} . The unit cell is shown with a dashed square. The blue atoms are the magnetic capping atoms. (b) Side view of the B_{11} . The B_I , B_{II} , and B_{III} are the three non-equivalent B atoms. (c) Angle view of the B_{11} . (d) Total energies as a function of the lattice constant for the FM and NM B_{11} , in which the minimum total energy for the FM state is set as energy zero. (e) Calculated phonon spectrum for the FM B_{11} . (f) Calculated energy evolution for the FM B_{11} at 133 K. The inset gives the snapshot of the final frame in the MD simulation.

TB model

One TB model is constructed to understand and generalize the magnetism.

$$H(k) = H_{on\ site}(k) + H_{hop}(k)$$

where $H_{on\ site}(k)$ and $H_{hop}(k)$ express the on site and hopping interactions, respectively. And they are written as

$$H_{on\ site}(k) = \begin{bmatrix} \varepsilon_0 & 0 \\ 0 & \varepsilon_0 \end{bmatrix}$$

$$H_{hop}(k) = \begin{bmatrix} T & t_0 \\ t_0 & T \end{bmatrix}$$

$$T = 2t_1(\cos k_x a + \cos k_y a)$$

where ε_0 represents the on-site energy of the top/bottom site, t_0 represents the strength of hopping interactions between the top and bottom sites mediated by the effective bridging atom in the structural unit, and t_1 represents the strength of hopping interactions between the sites in the nearest-neighboring unit cells. We can solve the equation to get the $E(k)$,

$$E(k) = \varepsilon_0 \pm t_0 + 2T$$

Magnetism mechanism

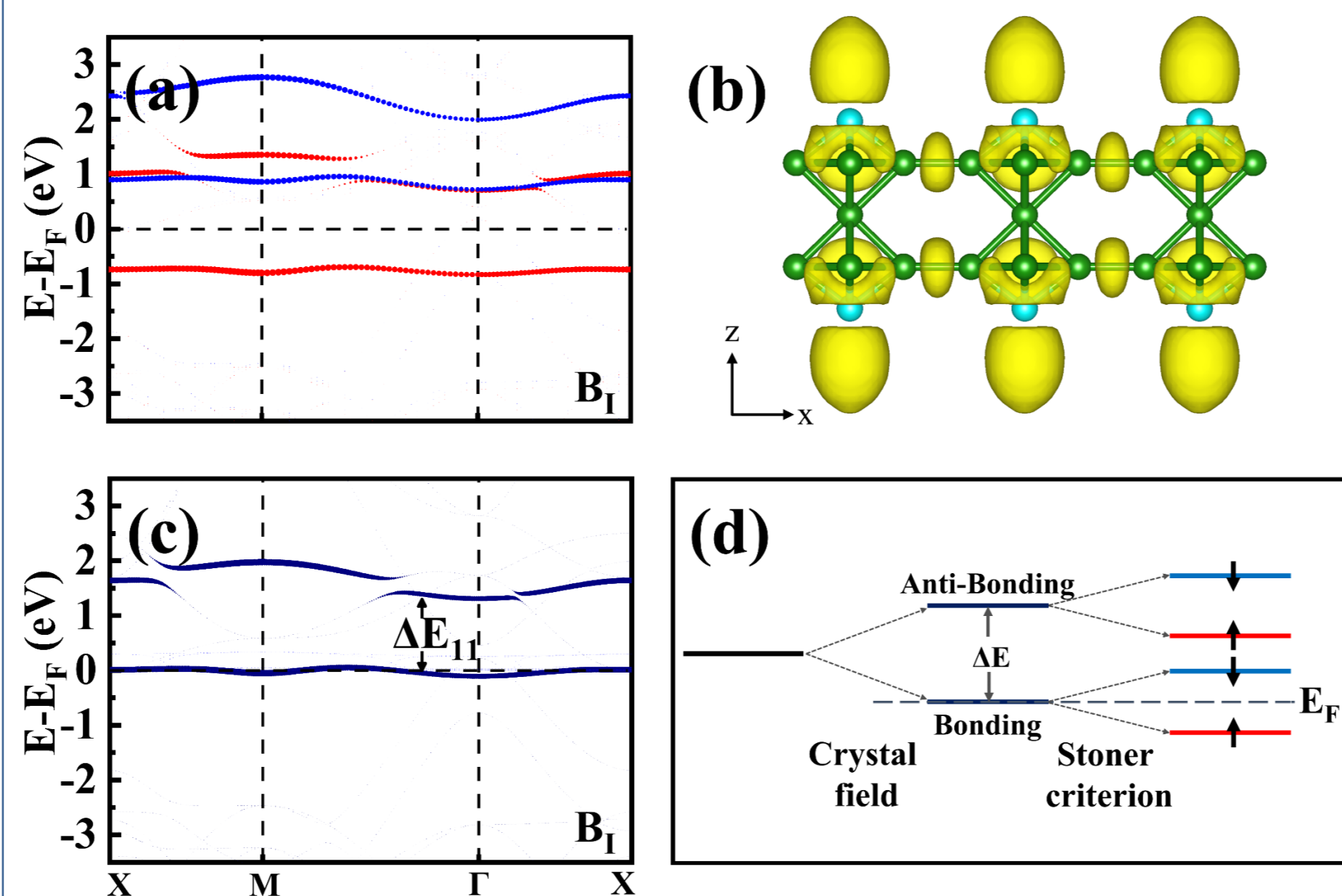


FIG. 3: (a) Projected bands of the $B_1 p_z$ orbital in the FM B_{11} . (b) Side view of the ELF of the FM borophene. (c) Projected bands of the $B_1 p_z$ orbital in the NM B_{11} . ΔE_{11} expresses the splitting strength of the $B_1 p_z$ bonding and anti-bonding states. (d) Magnetism mechanism diagram of the B_{11} .

Electronic Structures

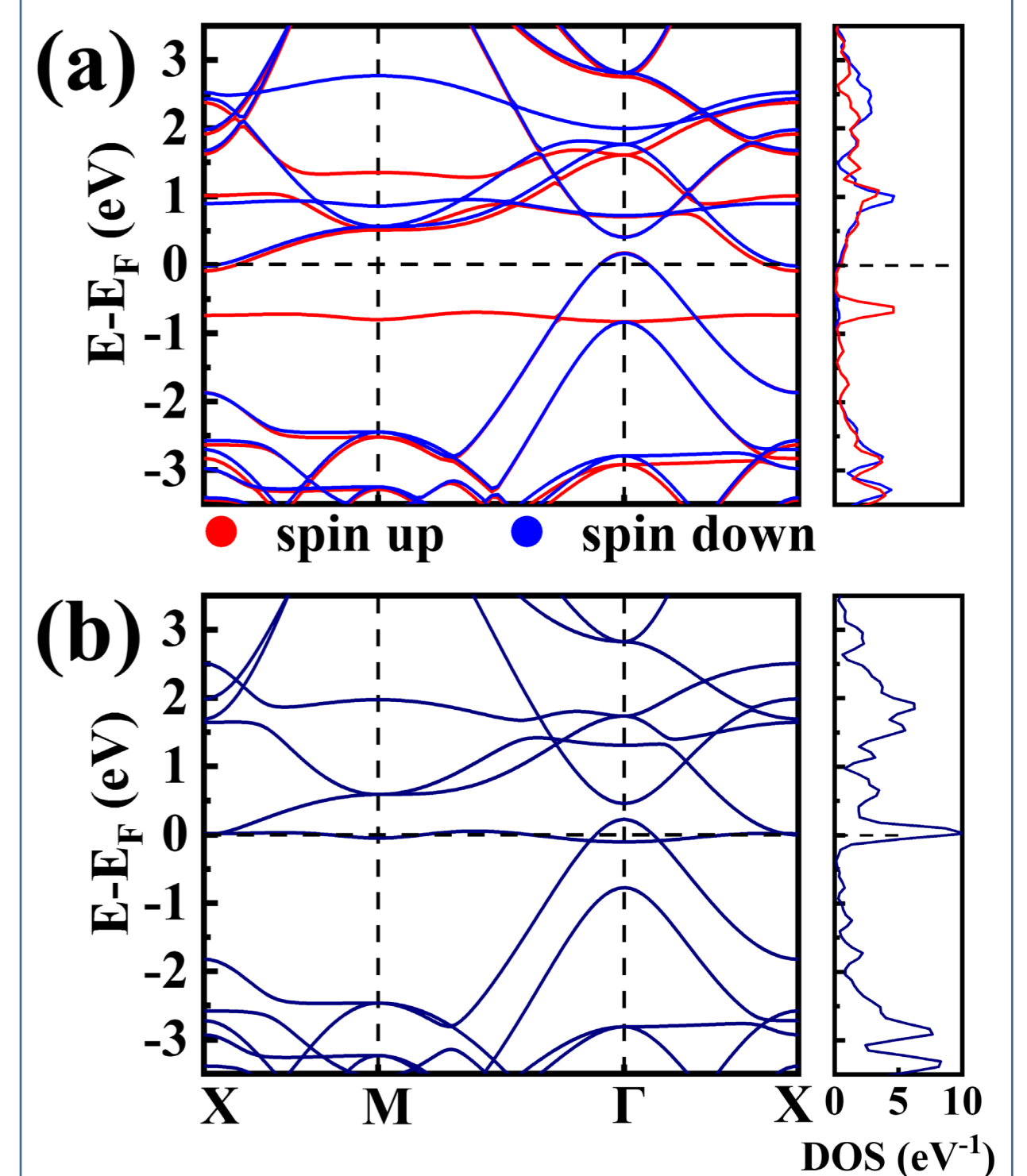


FIG. 2: (a) Band structures and DOSs for the magnetic B_{11} . The red and blue curves give the spin-up and spin-down components, respectively. (b) Band structure and DOS for the NM B_{11} .

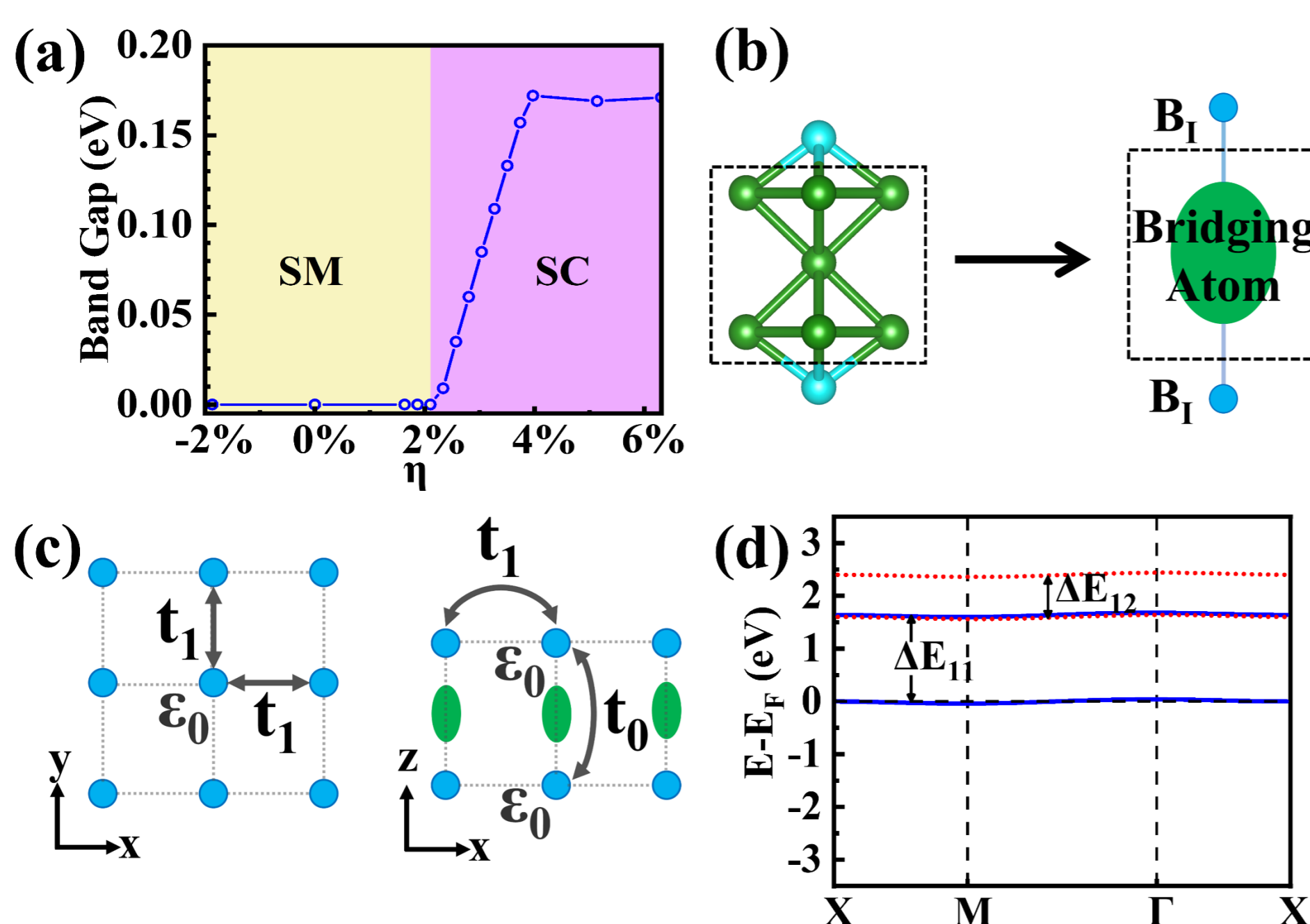


FIG. 4: (a) Band gap for the B_{11} under biaxial strain. The “SM” and “SC” represent the semimetal and semiconductor, respectively. (b) A general model proposed for the magnetism origin. (c) The simplified lattice adopted in the TB model. (d) The obtained TB bands for the B_{11} (blue curves) and B_{12} (red curves), respectively. The parameters for B_{11} are $\varepsilon_0 = 0.8$ eV, $t_0 = 0.8$ eV, and $t_1 = 0.01$ eV. The parameters for B_{12} are $\varepsilon_0 = 2.0$ eV, $t_0 = 0.4$ eV, and $t_1 = 0.01$ eV.

Conclusions

In summary, we have proposed a stable magnetic B_{11} monolayer, which has a tetragonal lattice and is a semimetal. The magnetism is triggered by the bonding flat band, formed due to the strong interactions between the p_z orbitals of the two capping B atoms mediated with middle B atoms in the structural unit. Under biaxial tensile strain, the B_{11} monolayer can transform from a semimetal to a semiconductor with a global band gap up to 171 meV. One tight binding model is constructed to explore and generalize the magnetism. One new borophene of nonmagnetic B_{12} is tuned to being magnetic based on the proposed tight binding model. Our findings expand the 2D magnetic materials and pave a new path for fabricating new 2D magnetic material composed of light elements.

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References

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