

Precise Raman Control of Spinor Matterwave with Nanosecond Composite Biased Rotations



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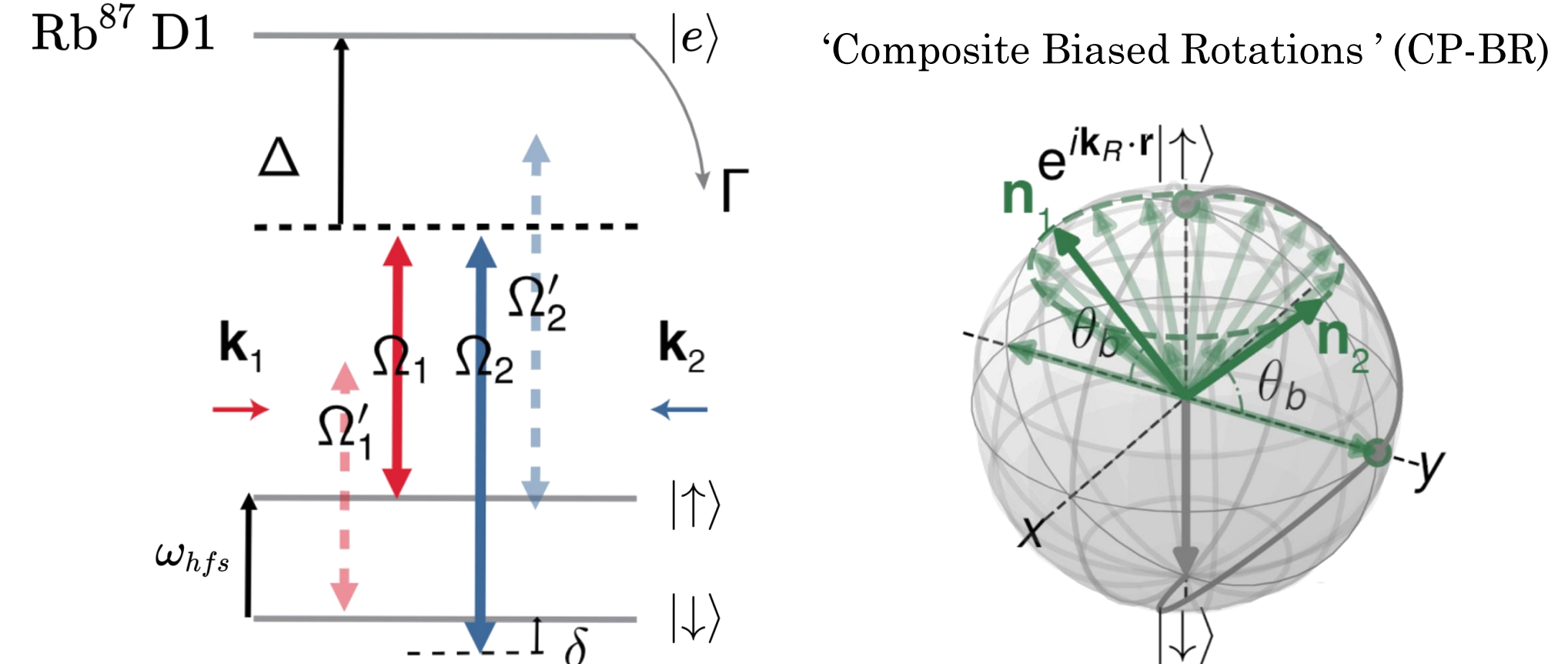
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I. Introduction

We experimentally demonstrate precise Raman matterwave control at an intermediate single-photon detuning $\Delta = O(\omega_{hfs})$, where a balance between the optical power efficiency with the requirements on the control speed and the suppression of excited-state dynamics can be adjusted. The method is based on composite biased rotation^[1] that exploits the proportionality between the traditionally “unwanted” light shift δ with the Raman coupling Ω_R . At $\Delta \approx 4\omega_{hfs}$, mesoscopic samples of 10^5 ⁸⁷Rb atoms are uniformly controlled, within **tens of nanoseconds**, near a laser focus with merely ~ 10 mW power. The control is fast enough to be immune to low-frequency noises, so our system can be accurately modeled. The $\mathcal{F} > 99.2\%$ fidelity is estimated with standard single-qubit QPT^[2] and RB^[3]. Our work suggests highly precise spinor matterwave controls are achievable for large atomic samples with moderate laser power, even in noisy environment.

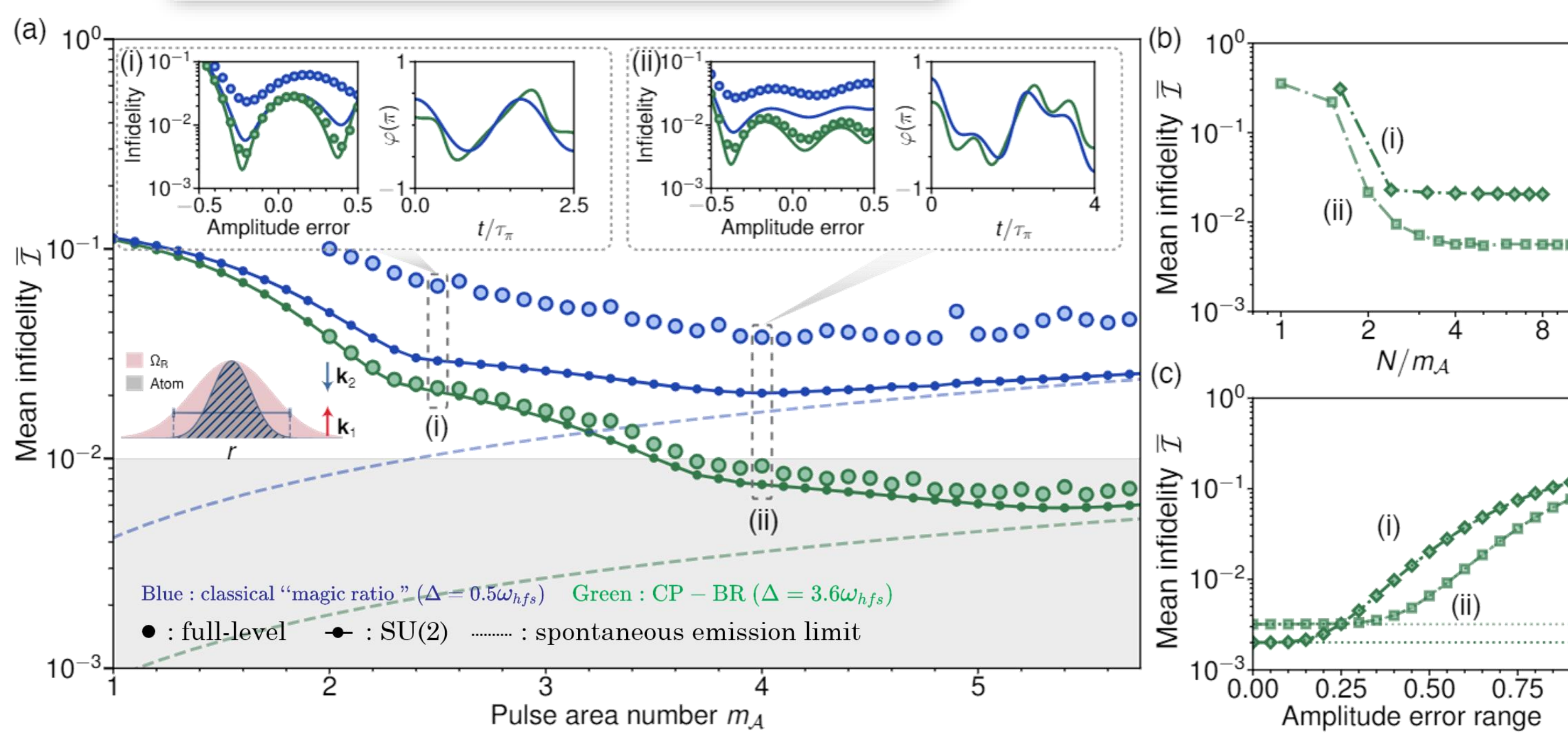


II. Theoretical Model

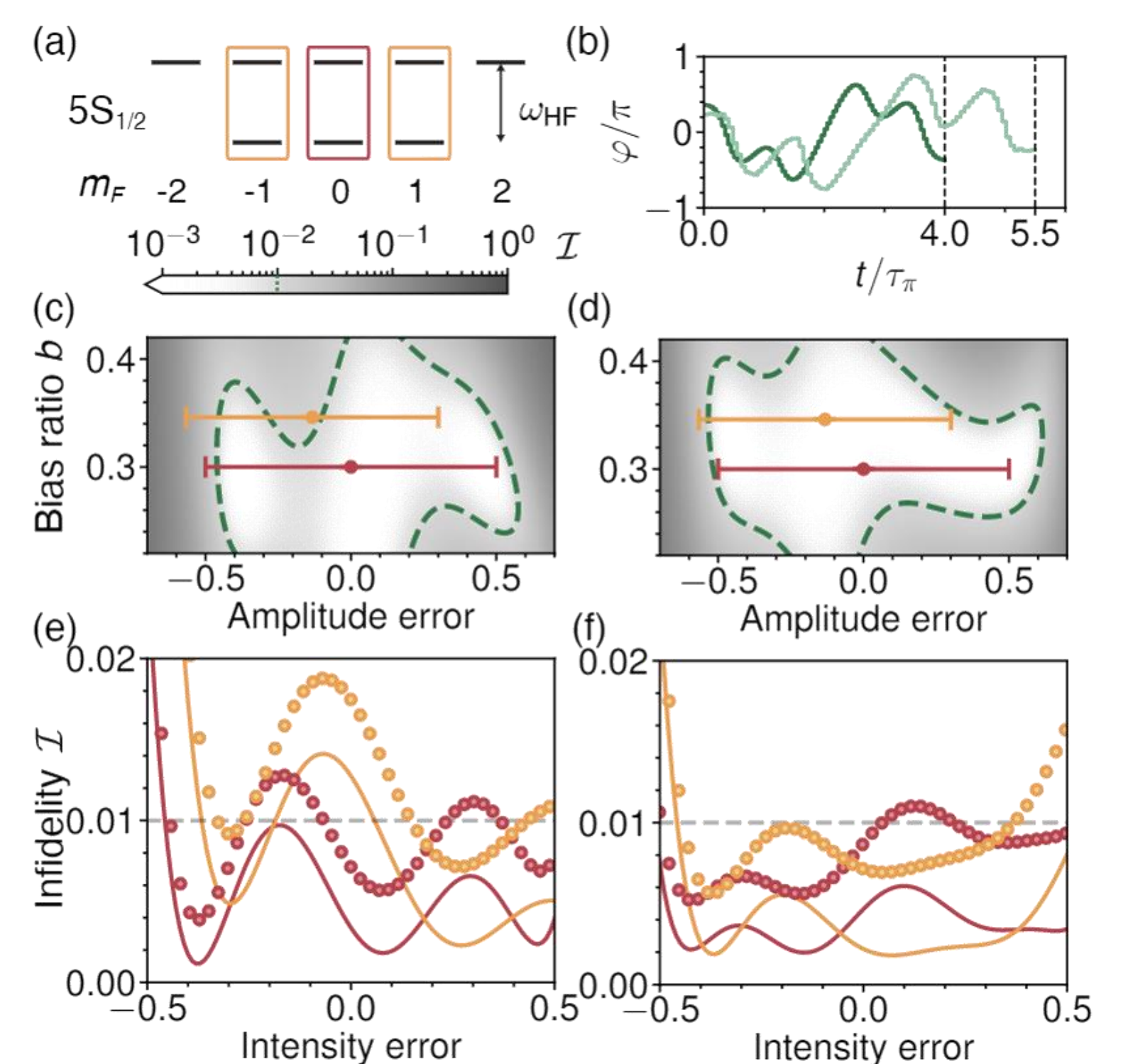
$$H = -\frac{1}{2}\hbar\delta\sigma_z + \frac{1}{2}\hbar|\Omega_R|(\sigma_x \cos\varphi + \sigma_y \sin\varphi)$$

$$\delta = \frac{|\Omega_1|^2}{4\Delta} + \frac{|\Omega_2|^2}{4(\Delta - \omega_{HF})} - \frac{|\Omega_2|^2}{4\Delta} - \frac{|\Omega_1|^2}{4(\Delta + \omega_{HF})}$$

$$\theta_b = \tan^{-1}(\delta/|\Omega_R|)$$



• Multiple Zeeman-spins



Intensity-error immunity supports high-fidelity matterwave control of multiple Zeeman spinors in parallel.

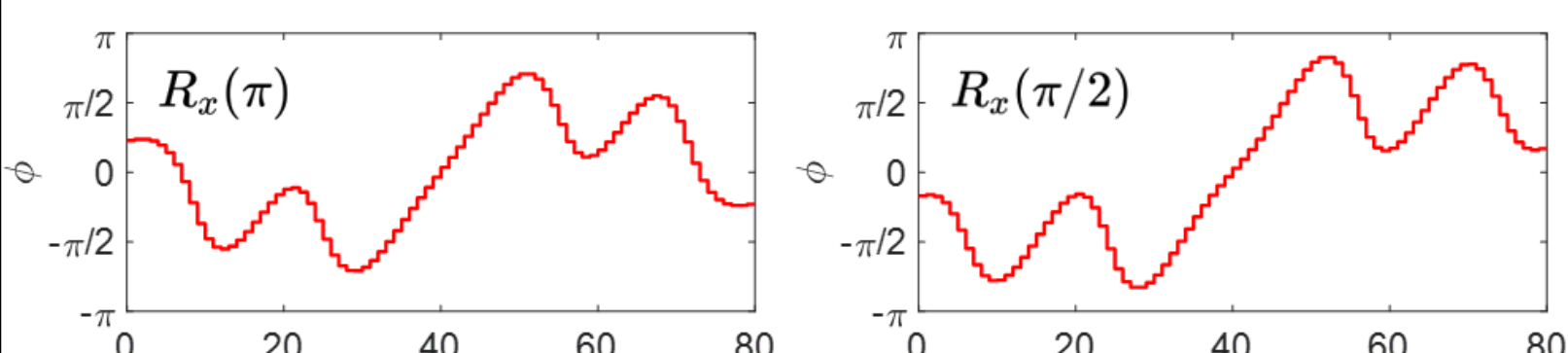
III. Pulse shaping

• Optimization algorithm-GRAPE

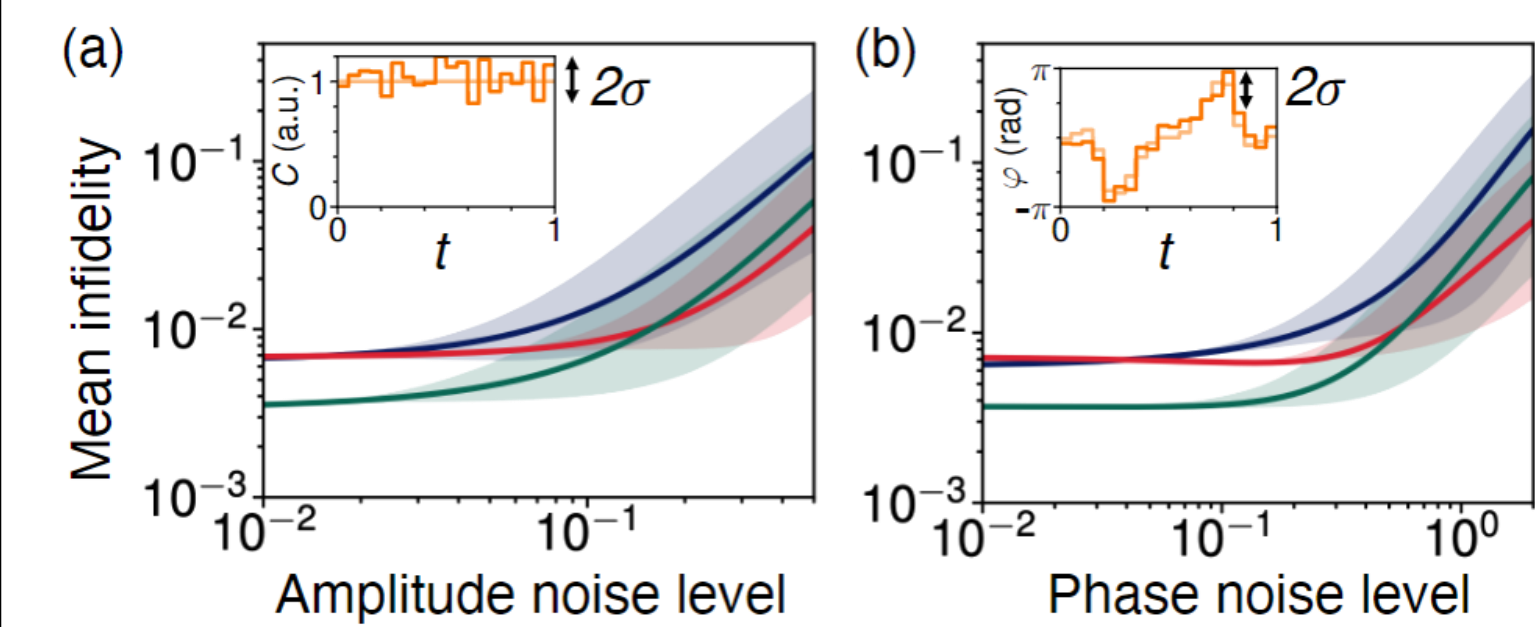
SU(2) rotation operator: $U(A, b, \varphi) = 1 \cos \frac{\varphi}{2} - i \sin \frac{\varphi}{2} \frac{\sigma_x \cos \varphi + \sigma_y \sin \varphi + b\sigma_z}{\sqrt{1+b^2}}$

Gate fidelity: $\mathcal{F} = \frac{1}{6} \sum_{j=1}^6 |\langle \psi_j | U^\dagger \tilde{U} | \psi_j \rangle|^2$

Gradient: $g_i = -\frac{\partial \mathcal{F}}{\partial \tilde{U}} \frac{\partial \tilde{U}}{\partial \varphi_i}$, $\frac{\partial \tilde{U}}{\partial \varphi_i} = \tilde{U}_N \dots \tilde{U}_{i+1} \frac{\partial \tilde{U}_i}{\partial \varphi_i} \tilde{U}_{i-1} \dots \tilde{U}_1$

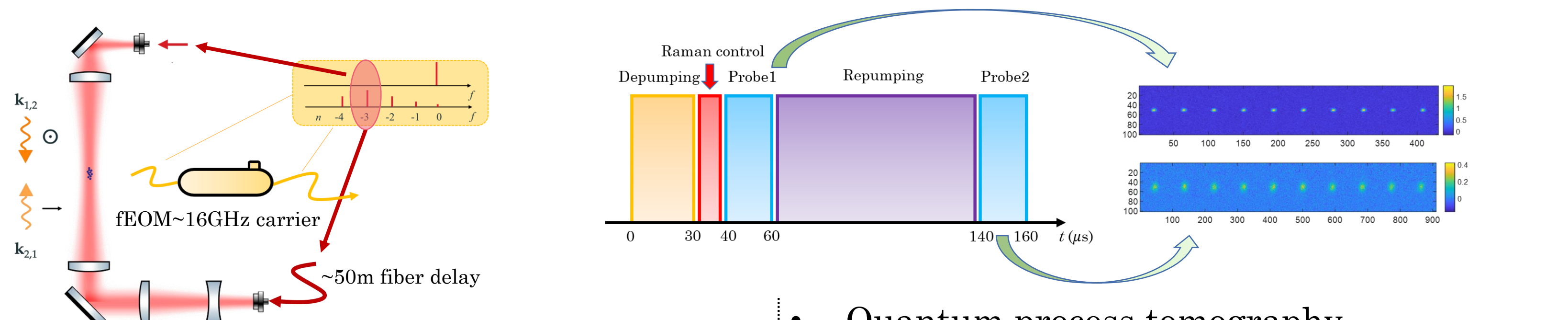


• Pulse noise resilience

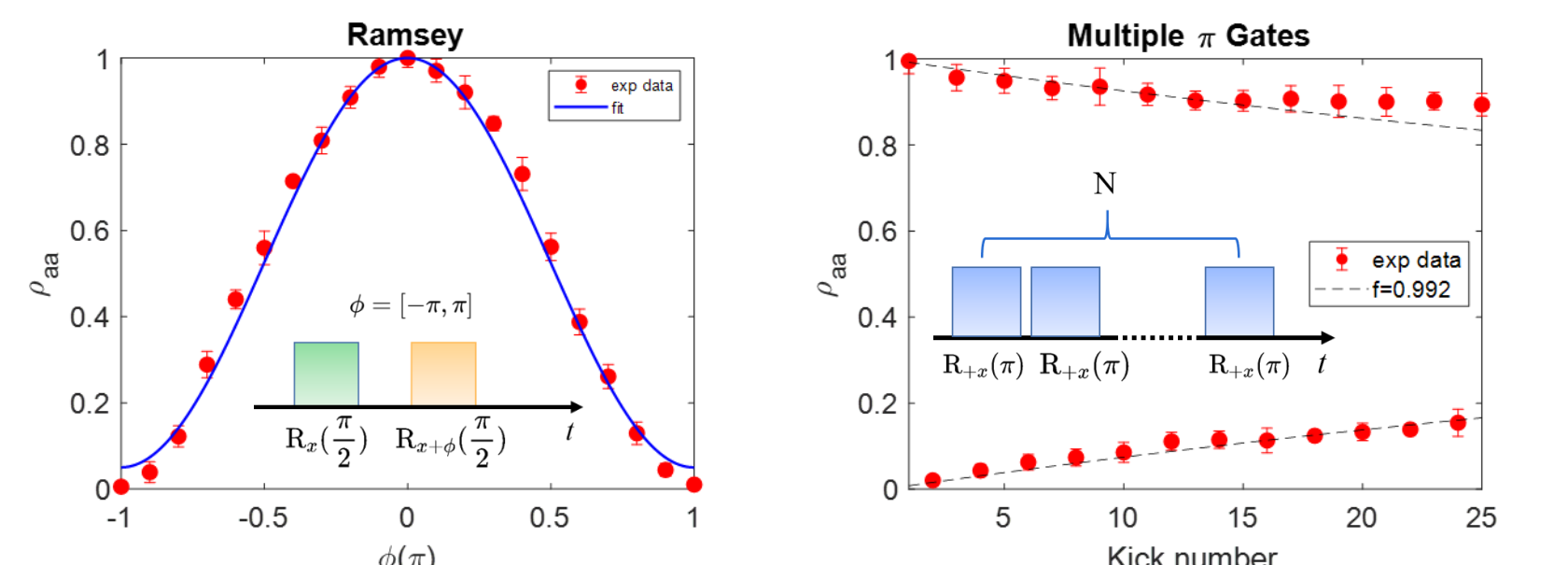


IV. Experimental setup and results

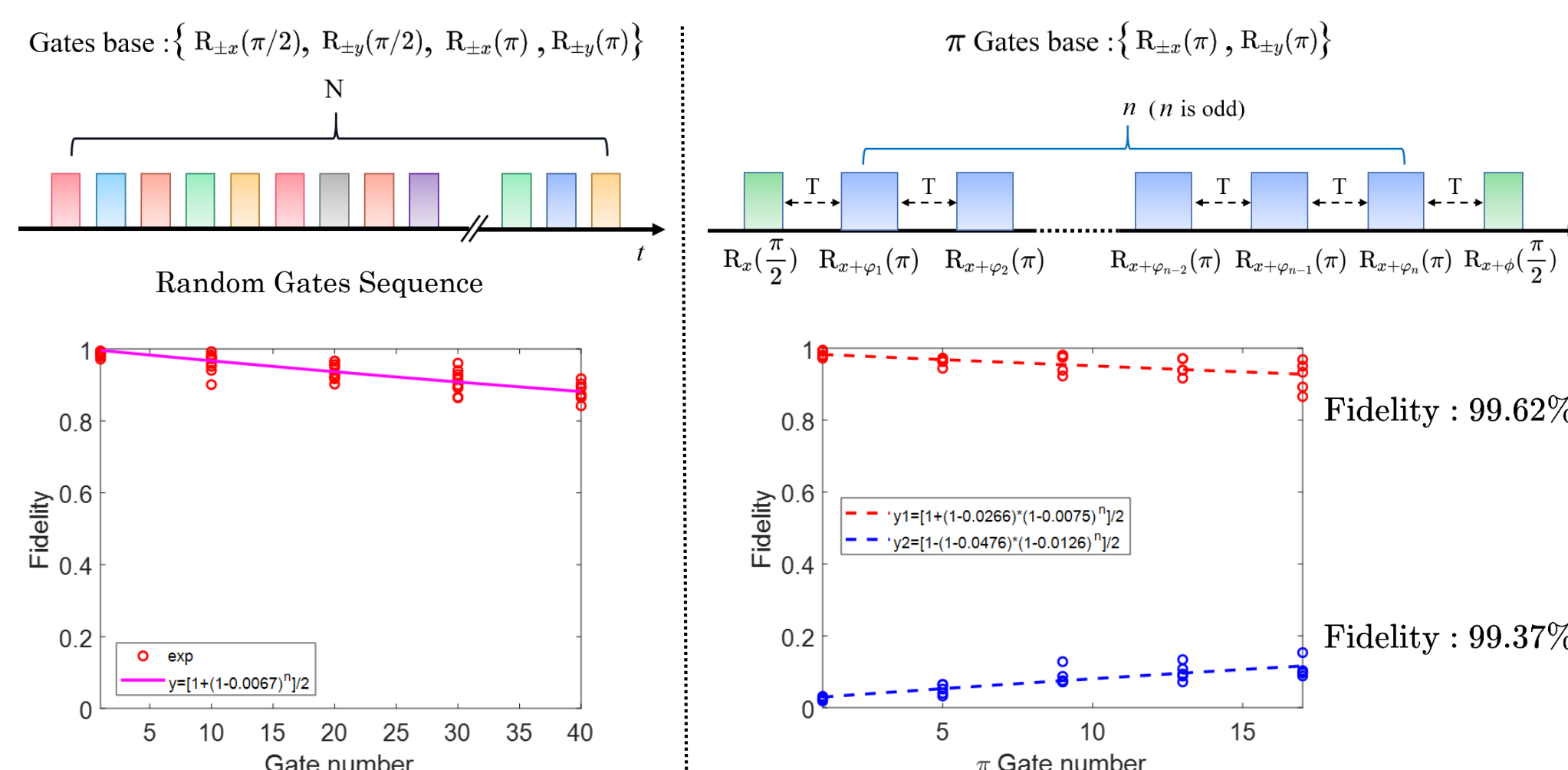
• The experimental setup and measurement procedure



• Ramsey Interferometry and Multiple π -Gates



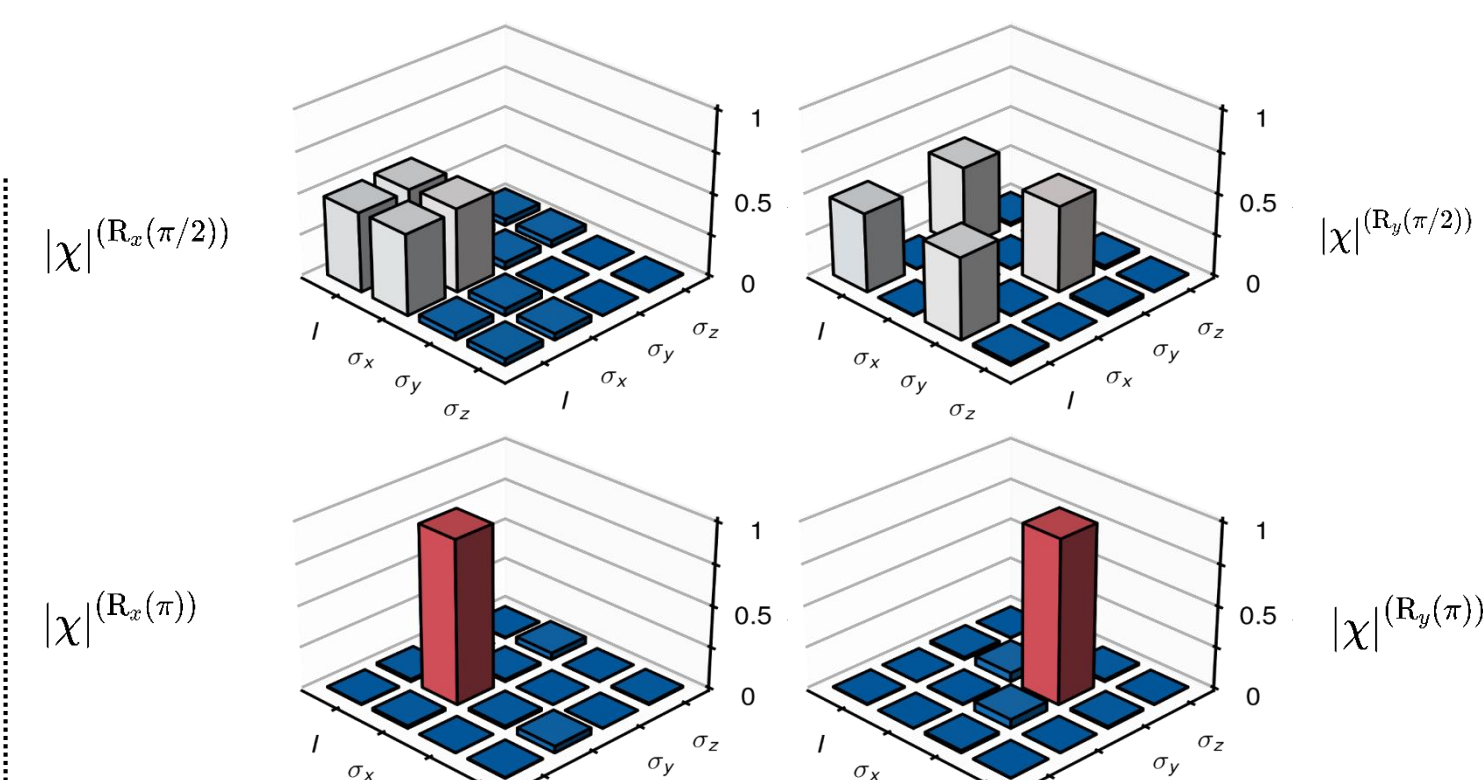
• Randomized benchmarking and dynamic decoupling^[2,3]



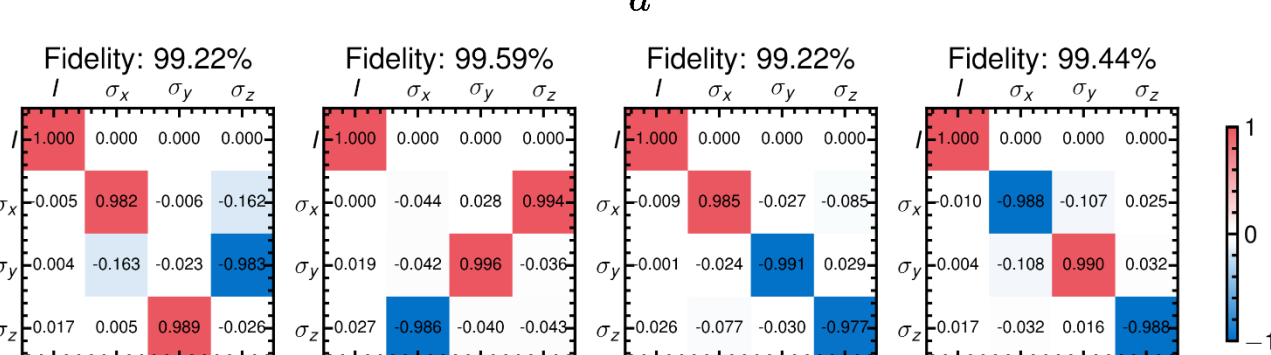
• Quantum process tomography

Gates base: $\{R_{+x}(\pi/2), R_{+y}(\pi/2), R_{+x}(\pi), R_{+y}(\pi)\}$

Gate χ -matrix: $\Lambda(\rho) = \sum_{j,k=1}^d \chi_{jk} P_j \rho P_k$, $P_j \in \mathcal{P}^{\otimes n}$, $\mathcal{P} = \{I, \sigma_x, \sigma_y, \sigma_z\}$



Pauli transfer matrix: $(R_A)_{ij} = \frac{1}{d} \text{Tr}\{P_i \Lambda(P_j)\}$



V. Discussions

- The 99.2% fidelity is less than the $\mathcal{F} = 99.7\%$ theoretical limit, likely due to digital noise in the waveforms (to be suppressed with new technology).
- Doubling Δ and laser power together lead to $\mathcal{F} \sim 99.9\%$.
- Toward a practical large area atom interferometer for wideband, noise-immune inertial sensing.
- Scaling-up the Raman control for advanced quantum information processors.

VI. Reference

- [1] QIU L, YUAN H, WU S. PHYSICAL REVIEW RESEARCH 5, 043094 (2023) High-fidelity Raman matterwave control by composite biased rotations[J].
 [2] GREENBAUM D. Introduction to Quantum Gate Set Tomography[J]. Quantum Physics, arXiv, 2015.
 [3] KNILL E, LEIBFRIED D, REICHLER R, et al. Randomized Benchmarking of Quantum Gates[J].